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Phase Transition in a Vlasov-Boltzmann Binary Mixture

Abstract There are not many kinetic models where it is possible to prove bifurcation phenomena for any value of the Knudsen number. Here we consider a binary mixture over a line with collisions and long range repulsive interaction between different species. It undergoes a segregation phase transition at sufficiently low temperature. The spatially homogeneous Maxwellian equilibrium corresponding to the mixed phase, minimizing the free energy at high temperature, changes into a maximizer when the temperature goes below a critical value, while non homogeneous minimizers, corresponding to coexisting segregated phases, arise. We prove that they are dynamically stable with respect to the Vlasov-Boltzmann evolution, while the homogeneous equilibrium becomes dynamically unstable.

1 Introduction and Main Results.

The phenomenon of phase transition in a thermodynamic system is usually described by the arising of multiple minimizers of the free energy. Namely, when the temperature is lowered below a certain critical value, the unique equilibrium minimizer becomes a local maximizer and new minimizers appear. This is interpreted as a loss of stability of the old minimizer and the

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birth of new stable states. Here the meaning of stability is merely related to the free energy minimizing properties of the state. Of course it would be desirable to have a more detailed, dynamical analysis of the stability properties involved in such a phenomenon of phase transition. This issue is, generally speaking, very difficult, because the natural dynamics is a many-component one, whose detailed understanding is very far from being attainable. A simpler but, in our opinion, still interesting problem is to study this kind of behavior at kinetic level, where, instead of a system of a huge number of particles, one has to study a partial differential equation for the probability distribution of particles on the one particle phase space. This is the problem we want to address in this paper.

The standard kinetic model of a rarefied gas undergoing collisions (Boltzmann equation), however, describes essentially the ideal gas, which does not exhibit phase transitions. Following van der Waals, it is convenient to include a long range attractive interaction between particles to see a vapor-liquid transition. With the introduction of such an interaction, new problems arise, due to the fact that nothing prevents the system from collapsing (Statistical Mechanics instability). The van der Waals approach of adding a hard core interaction is not easy to handle and more complicated, many body long range interactions have been introduced (see [13,15]) to handle this in the framework of the equilibrium Statistical Mechanics.

While one could in principle try to follow the approach in [13], we find easier to consider a different kinetic model introduced a few years ago in [2], where there is no attractive interaction and the Statistical Mechanics instability does not arise. The model consists of two species of particles which, for simplicity, have the same mass. To fix the ideas, think of them as distinguished just by their *color*, say red and blue. Their short range interaction is modeled by Boltzmann-like collisions which are color blind, while the long range repulsive interaction, arising only between particles of different color, is modeled by a Vlasov force with a smooth, bounded, finite range potential. We refer to [2], [1],[3], [4] for more information on this model. For any $t \in \mathbb{R}^+$, let the non negative normalized functions $f_i(t, x, \xi)$, $i = 1, 2$, denote the probability densities of finding a particle of species 1 (red) or 2 (blue) in a cell of the phase space around the point (x, ξ) at time t . In this paper we will only consider the case $x \in \mathbb{R}$, the real line, while the velocity $\xi = (v, \zeta) \in \mathbb{R}^3$ with $\zeta \in \mathbb{R}^2$. The time evolution is governed by

$$\partial_t f_1 + v \partial_x f_1 + F(f_2) \partial_v f_1 = Q(f_1, f_1 + f_2), \quad (1.1)$$

$$\partial_t f_2 + v \partial_x f_2 + F(f_1) \partial_v f_2 = Q(f_2, f_1 + f_2).$$

Here the Vlasov force $F(h)$ due to the mass distribution h is defined as

$$F(h)(t, x) = -\partial_x \int_{\mathbb{R}} dy U(|x-y|) \int_{\mathbb{R}^3} d\xi h(t, y, \xi), \quad (1.2)$$

where $U(r)$, the interaction potential is non negative, smooth, bounded, with $U(r) = 0$ for $r \geq 1$ and

$$\int_{\mathbb{R}} dx U(|x|) = 1.$$

The collision integral is defined as:

$$\begin{aligned} Q(h_1, h_2)(\xi) &= \int_{\mathbb{R}^3 \times \mathbb{S}^2} |(\xi - \xi') \cdot \omega| \{h_1(\xi_*)h_2(\xi'_*) - h_1(\xi)h_2(\xi')\} d\xi' d\omega \\ &:= Q_{\text{gain}}(h_1, h_2) - Q_{\text{loss}}(h_1, h_2), \end{aligned} \quad (1.3)$$

with $\mathbb{S}^2 = \{\omega \in \mathbb{R}^3 \mid |\omega| = 1\}$ and ξ_*, ξ'_* related to ξ, ξ' by the usual elastic collision relations

$$\xi_* = \xi - \omega[\omega \cdot (\xi - \xi')], \quad \xi'_* = \xi' + \omega[\omega \cdot (\xi - \xi')]. \quad (1.4)$$

For any $\beta > 0$, the couple $(a_1\mu_\beta, a_2\mu_\beta)$, with μ_β the spatially homogeneous Maxwellian at temperature β^{-1} ,

$$\mu_\beta(\xi) = \left(\frac{\beta}{2\pi}\right)^{\frac{3}{2}} e^{-\beta\xi^2/2} \quad (1.5)$$

and $a_i > 0$, is an equilibrium solution and the most general homogeneous equilibrium differs from it just for rescaling and centering. However, due to the presence of the Vlasov force, non homogeneous Maxwellian equilibria are possible and they are of the form

$$f_i(x, \xi) = \rho_i(x)\mu_\beta(\xi) \quad (1.6)$$

provided that the densities $\rho_i(x) > 0$ satisfy the conditions

$$\log \rho_i + \beta U * \rho_{i+1} = C_i, \quad i = 1, 2 \quad (1.7)$$

Here and in the rest of the paper the label $i + 1$ means 2 if $i = 1$ and 1 if $i = 2$ and the convolution product $*$ is defined by $a * b(x) = \int_{\mathbb{R}} dy a(|x - y|)b(y)$. The constants C_i have the physical meaning of chemical potentials. The phase transition phenomenon we want to discuss occurs when $C_1 = C_2$, a situation where the equilibrium solutions are symmetric under the exchange $1 \leftrightarrow 2$. This symmetry is spontaneously broken below the critical temperature. Therefore we assume $C_1 = C_2$ in the rest of the paper.

If β is suitably small, conditions (1.7) are satisfied only by constant ρ_i 's. On the other hand if β is sufficiently large, non constant ρ_i solving (1.7) can be constructed.

To understand the arising of multiple equilibria, let us start by defining the local free energy $\varphi(\rho_1, \rho_2)$ on $\mathbb{R}^+ \times \mathbb{R}^+$ (i.e. the free energy density when U is replaced by a δ -function) as

$$\varphi(\rho_1, \rho_2) = \rho_1 \log \rho_1 + \rho_2 \log \rho_2 + \beta \rho_1 \rho_2 \quad (1.8)$$

Let $\rho = \rho_1 + \rho_2$. It is shown in [4] (in a slightly more general context) that, if $\beta\rho < 2$ then the only stationary points for φ are characterized by $\rho_1 = \rho_2 = \rho/2$ and they are minimizers for φ . On the other hand, if $\beta\rho > 2$, then there are $\rho^+ > \rho^- > 0$ such that

$$(\rho_1, \rho_2) = (\rho^+, \rho^-) \quad (1.9)$$

(and $(\rho_1, \rho_2) = (\rho^-, \rho^+)$ by symmetry) is an absolute minimizer for φ , while $(\rho_1, \rho_2) = (\rho/2, \rho/2)$ is a local maximizer for φ . When (a_1, a_2) is chosen equal to (ρ^+, ρ^-) , the equilibrium state is interpreted as a pure phase rich of red particles, while the state $(a_1, a_2) = (\rho^-, \rho^+)$ corresponds to a phase rich of blue particles.

Non homogeneous solutions become relevant when one has to describe a situation of phase coexistence. The idea to construct non homogeneous solutions is to observe that at low temperature, in order to minimize the local free energy, the system has to be in a pure phase at infinity. A situation with phase coexistence should arise when the system is in the blue-rich phase at $-\infty$ and in the red-rich phase at $+\infty$ or vice versa. Therefore one looks at the minimizers of a suitable free energy functional with above constraints. Indeed, the conditions (1.7) are the Euler-Lagrange equations for the *free energy* functional

$$\mathcal{F}(\rho_1, \rho_2) = \int_I dx (\rho_1 \log \rho_1 + \rho_2 \log \rho_2) + \beta \int_I dx \int_I dy U(|x-y|) \rho_1(x) \rho_2(y). \quad (1.10)$$

on a finite interval I with periodic boundary conditions. Moreover, it can be shown, (see [3]), that when β is large, with the masses of the two species fixed, there are non constant couples (ρ_1, ρ_2) giving lower values to the free energy than the constants. Over the real line $-\infty < x < \infty$, above free energy does not make sense and a more careful definition is required: Let $(-\ell, \ell)$ be a bounded interval and set

$$\begin{aligned} \mathcal{F}_{(-\ell, \ell)}(\rho_1, \rho_2) &= \int_{(-\ell, \ell)} dx \varphi(\rho_1, \rho_2) \\ &+ \frac{\beta}{2} \int_{(-\ell, \ell)^2} dx dy U(|x-y|) [\rho_1(x) - \rho_1(y)] [\rho_2(y) - \rho_2(x)]. \end{aligned}$$

We define the *excess free energy* as

$$\hat{\mathcal{F}}(\rho_1, \rho_2) := \lim_{\ell \rightarrow \infty} [\mathcal{F}_{(-\ell, \ell)}(\rho_1, \rho_2) - 2\ell \varphi(\rho^+, \rho^-)]. \quad (1.11)$$

Note that, since $\varphi(\rho^+, \rho^-) = \varphi(\rho^-, \rho^+)$, the functional $\hat{\mathcal{F}}$ is $+\infty$ when $\rho = (\rho_1, \rho_2)$ does not go to pure phases at infinity. Since we are interested at the coexistence of phases, we assume that (ρ_1, ρ_2) satisfy one of the conditions $\lim_{x \rightarrow \pm\infty} \rho_1(x) = \rho^\pm$ or $\lim_{x \rightarrow \pm\infty} \rho_2(x) = \rho^\mp$. In [4] several results are proved about the minimizers for the functional $\hat{\mathcal{F}}$, which are summarized in the following theorem (see also [5]):

Theorem 1.1 *Let $\beta\rho > 2$. Then there exists a unique (up to translations) positive minimizer (front) for the one-dimensional excess free energy $\hat{\mathcal{F}}$, defined by (1.11), in the class of continuous front functions $\rho = (\rho_1(x), \rho_2(x))$ such that $\lim_{z \rightarrow \pm\infty} \rho_1 = \rho^\pm$, $\lim_{z \rightarrow \pm\infty} \rho_2 = \rho^\mp$, where ρ^\pm are defined in (1.9). We denote by $\bar{\rho} = (\bar{\rho}_1(x), \bar{\rho}_2(x))$ the unique minimizer such that $\bar{\rho}_1(x) = \bar{\rho}_2(-x)$. $\bar{\rho}_1$ is monotone increasing and $\bar{\rho}_2$ is monotone decreasing and $\rho^- < \bar{\rho}_i(x) < \rho^+$ for any $x \in \mathbb{R}$. Moreover, the front $\bar{\rho}$ is $C^\infty(\mathbb{R})$ -smooth*

and satisfies the Euler-Lagrange equations (1.7); Its derivative $\bar{\rho}'$ satisfies the equations

$$\bar{\rho}'_i + \beta \bar{\rho}_i (U * \bar{\rho}'_{i+1}) = 0, \quad i = 1, 2. \quad (1.12)$$

The front $\bar{\rho}$ converges to its asymptotic values exponentially fast, in the sense that there is $\alpha > 0$ such that

$$|\bar{\rho}_1(x) - \rho^\mp| e^{\alpha|x|} \rightarrow 0 \text{ as } x \rightarrow \mp\infty, \quad |\bar{\rho}_2(x) - \rho^\pm| e^{\alpha|x|} \rightarrow 0 \text{ as } x \rightarrow \mp\infty.$$

Finally, the derivatives of $\bar{\rho}_i$ of any order vanish at infinity exponentially fast and $\bar{\rho}'$ is odd in the sense that

$$\bar{\rho}'_1(x) = -\bar{\rho}'_2(-x). \quad (1.13)$$

The aim of this paper is to show that for β larger than a certain critical value, the non homogeneous equilibria, the front solutions $(\bar{\rho}_1, \bar{\rho}_2)$ are dynamically stable with respect to the evolution (1.1), while the homogeneous mixed phase is unstable. This will show that a complete bifurcation scenario arises in this model for any value of the Knudsen number, i.e. the ratio between the mean free path and the range of the interaction potential.

In the rest of the paper, without loss of generality we fix the asymptotic total density $\rho = \rho^+ + \rho^- = 2$. If $\beta < 1$, the only minimizer is the homogeneous equilibrium (mixed phase) $M_{\text{hom}}(\xi) = (\mu_\beta(\xi), \mu_\beta(\xi))$. On the other hand, if $\beta > 1$, the pure phases $M_{\text{red}} = (\rho^+ \mu_\beta(\xi), \rho^- \mu_\beta(\xi))$ and $M_{\text{blue}} = (\rho^- \mu_\beta(\xi), \rho^+ \mu_\beta(\xi))$ are constant minimizers and M_{hom} is a maximizer. Moreover, we have the non homogeneous equilibrium

$$M_{\bar{\rho}} = (\bar{\rho}_1(x) \mu_\beta(\xi), \bar{\rho}_2(x) \mu_\beta(\xi)), \quad (1.14)$$

with $\bar{\rho}_1(x) + \bar{\rho}_2(x) \rightarrow 2$ as $x \rightarrow \pm\infty$.

We assume that the initial datum for the dynamics (1.1) is a small perturbation of one of the above equilibria, denoted generically by M :

$$f_i(0) = M_i + \sqrt{M_i} g_i(0), \quad (1.15)$$

with g_i sufficiently small in a sense that will be specified later.

Let $f(t, x, \xi)$ be the solution to (1.1) and define $g(t, x, \xi)$ by setting

$$f_i(t, x, \xi) = M_i(x, \xi) + \sqrt{M_i(x, \xi)} g_i(x, \xi, t). \quad (1.16)$$

The equation for the perturbation g is

$$\begin{aligned} & \left\{ \partial_t + v \partial_x + F \left(M_{i+1} + \sqrt{M_{i+1}} g_{i+1} \right) \partial_v + \nu(x, \xi) \right\} g_i \\ &= \beta F \left(\sqrt{M_{i+1}} g_{i+1} \right) \sqrt{M_i} v + \sum_{j=1}^2 K^{i,j} g_j + F \left(\sqrt{M_{i+1}} g_{i+1} \right) g_i v \\ &+ \Gamma(g_i, g_i) + \Gamma(g_i, g_{i+1}), \end{aligned} \quad (1.17)$$

where we have used the notation:

$$\begin{aligned}
\nu(x, \xi) &= \int_{\mathbb{R}^3 \times \mathbf{S}^2} |(\xi - \xi') \cdot \omega| (M_1(\xi') + M_2(\xi')) d\xi' d\omega, \\
K^{i,i} g_i &= \frac{1}{\sqrt{M_i}} Q_{\text{gain}} \left(\sqrt{M_i} g_i, M_1 + M_2 \right) + \frac{1}{\sqrt{M_i}} Q \left(M_i, \sqrt{M_i} g_i \right), \\
K^{i,i+1} g_{i+1} &= \frac{1}{\sqrt{M_i}} Q \left(M_i, \sqrt{M_{i+1}} g_{i+1} \right), \\
\Gamma(g_i, g_j) &= \frac{1}{\sqrt{M_i}} Q \left(\sqrt{M_i} g_i, \sqrt{M_j} g_j \right).
\end{aligned} \tag{1.18}$$

We will need the following symmetry condition on the initial data:

$$g_1(0, x, v, \zeta) = g_2(0, -x, -v, \zeta). \tag{1.19}$$

We note that such a property is preserved by the time evolution. It plays a crucial role in removing the obvious orbital instability of our problem related to the translation invariance: indeed such an invariance is explicitly broken when condition (1.19) is assumed.

The momentum, kinetic energy and particle number of each species are conserved during collisions, while the sum of potential and kinetic energy is conserved along the trajectories. Therefore the following quantities are conserved under the evolution (1.1):

- The masses of the perturbation g ,

$$\mathcal{M}_i(g) = \int_{\mathbb{R}^3} dx \int_{\mathbb{R}^3} d\xi (f_i(x, \xi) - M_i(x, \xi)), \quad i = 1, 2;$$

- The energy of the perturbation g ,

$$\begin{aligned}
\mathcal{E}(g) &= \int_{\mathbb{R}^3} dx \left[\int_{\mathbb{R}^3} d\xi \frac{\xi^2}{2} [(f_1 + f_2) - (M_1 + M_2)] \right. \\
&\quad \left. + \int_{\mathbb{R}^3} dy U(|x - y|) [\rho_{f_1}(x) \rho_{f_2}(y) - \rho_{M_1}(x) \rho_{M_2}(y)] \right], \tag{1.20}
\end{aligned}$$

with $f_i = M_i + \sqrt{M_i} g_i$ and

$$\rho_f(x) = \int_{\mathbb{R}^3} d\xi f(x, \xi), \tag{1.21}$$

the spatial density of the distribution f .

Moreover, by standard arguments it follows that the H -function of the perturbation g ,

$$\begin{aligned}
H(g) &= \int_{\mathbb{R}^3} dx \int_{\mathbb{R}^3} d\xi \left[(f_1(x, \xi) \log f_1(x, \xi) + f_2(x, \xi) \log f_2(x, \xi)) \right. \\
&\quad \left. - (M_1(x, \xi) \log M_1(x, \xi) + M_2(x, \xi) \log M_2(x, \xi)) \right] \tag{1.22}
\end{aligned}$$

does not increase in the evolution (1.1).

We notice that, since the spatial domain is \mathbb{R} , the total masses, energy and H -function of the distribution f are not well defined, while above differences are finite. Since $H(g)$ does not increase during the evolution and $\mathcal{E}(g)$ and the total masses $\mathcal{M}_i(g)$ are constant, any linear combination of them, with positive coefficient for the entropy, does not increase. In particular, the following non-increasing *entropy-energy* functional is crucial to study the stability of the equilibria:

$$\mathcal{H}(g) = H(g) + \beta \mathcal{E}(g) - \sum_{i=1}^2 \mathcal{M}_i(g) \left\{ C_i + 1 + \ln \left(\frac{\beta}{2\pi} \right)^{\frac{3}{2}} \right\}. \quad (1.23)$$

The factors multiplying the masses are suitably chosen to cancel some linear terms, as will be shown in next section. The factor β in front of the temperature is dictated by the free energy minimizing properties of the equilibria.

In the next sections we shall use the following weighted norm for $p \in [1, +\infty]$:

$$\|f\|_{L_w^p} = \left(\int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi |w(\xi) f(x, \xi)|^p \right)^{\frac{1}{p}},$$

for some positive weight function w . Moreover $L_{w,p}$ denotes the space of the measurable functions on $\mathbb{R} \times \mathbb{R}^3$ with $\|\cdot\|_{L_{w,p}}$ finite. When notation will be unambiguous, we will also omit the index w . Finally $\nabla_{x,v}$ denotes the couple (∂_x, ∂_v) .

Theorem 1.2 (Stability) *Assume $\beta > 1$ and $M = M_{\bar{\rho}}$. Let $w(\xi) = (\Sigma + |\xi|^2)^\gamma$, for $\Sigma > 0$ and $\gamma > 3/2$. There are $\delta > 0$ and $\Sigma > 0$, $C > 0$ such that, if the initial datum $f_i(0) = M_i + \sqrt{M_i} g_i(0)$ satisfies the symmetry condition (1.19), the bound*

$$\|wg(0)\|_{L^\infty} + \sqrt{\mathcal{H}(g(0))} < \delta$$

and $\|\nabla_{x,v}g(0)\|_{L^2} < +\infty$, then the initial value problem for (1.17) has a unique global in time solution with

$$\sup_{0 \leq t \leq \infty} \|wg(t)\|_{L^\infty} \leq C \{ \|wg(0)\|_{L^\infty} + \sqrt{\mathcal{H}(g(0))} \}, \quad (1.24)$$

$$\|\nabla_{x,v}g(t)\|_{L^2} \leq e^{Ct} \|\nabla_{x,v}g(0)\|_{L^2}. \quad (1.25)$$

Remark 1.1 Theorem 1.2 implies the stability of the non homogeneous equilibrium solution $M_{\bar{\rho}}$ in the Vlasov-Boltzmann evolution (1.1), with respect to initial perturbations satisfying (1.19). It turns out that $H(g(0)) > 0$ for $\|wg(0)\|_{L^\infty}$ small (see Lemma 2.2).

Dynamical stability and rate of convergence for the same front solutions has been established [5] for the Vlasov-Fokker-Planck dynamics: the Vlasov force is the same introduced here, but the collisions are replaced by a Fokker-Planck operator modeling the contact with a reservoir at inverse temperature β . A careful analysis of the *macroscopic* equation plays an important role.

To prove nonlinear stability, the first major mathematical difficulty we encounter is the presence of a large amplitude potential U . To our knowledge,

so far there has been no published work on stability result in the presence of a large external field in the Boltzmann theory. The main problem is the collapse of Sobolev estimate in higher order energy norms. Indeed, even upon taking one x -derivative, the H^1 -norm might actually grow in time due to the presence of the term $\partial_x F(M_{i+1})\partial_v g_i$ in (4.14). We are thus forced to design a strategy of proof based on a weighted L^∞ formulation without any derivatives and get control of derivatives only afterwards. Furthermore, unlike the previous linear Fokker-Planck interaction, the nonlinear Boltzmann collisions make it difficult to analyze the equation for the projection on the hydrodynamical modes in this case. In fact, even a L^2 stability of g is difficult to obtain directly from the analysis of eq. (1.17). Instead, we avoid a direct study of eq. (1.17) and make crucial use of the fundamental entropy-energy $\mathcal{H}(g)$ estimate to obtain a mixed $L^1 - L^2$ type of stability estimate, based on the spectral gap of the linearized free energy operator A (Lemma 2.2). We note that the spectral gap estimate relies in an essential way on the minimizing properties of the front solution and the symmetry condition is used to control the part of the solution in the null space of the operator A . We then bootstrap such a L^2 stability to a L^∞ estimate to obtain pointwise stability estimate, by following the curved trajectory induced by the force field $F(M_{i+1} + \sqrt{M_{i+1}}g_{i+1})$. The success of such a strategy is somewhat surprising: the perturbation of the field does not need to decay in time, and all the analysis over the one dimensional trajectory is carried out in a *finite time* interval $[0, T_0]$ (Lemma 4.1). Yet this is still sufficient for the global in time estimate due to the strong exponential time decay from the collision frequency.

Remark 1.2 The same result also holds when perturbing initially the equilibria M_{red} and M_{blue} . Indeed the proof is even simpler in this case, because the the analog of the operator A has again a spectral gap property, but its null space is trivial. Due to this, the symmetry assumption (1.19) is no more needed.

We next discuss the homogeneous equilibrium $M_{\text{hom}} = (\mu_\beta, \mu_\beta)$.

Theorem 1.3 *Assume $\beta < 1$ and $M = M_{\text{hom}}$. There are $\delta > 0$ and $\Sigma > 0$, $C > 0$ such that, if the initial datum $f_i(0) = M_i + \sqrt{M_i}g_i(0)$ satisfies the bound*

$$\|wg_0\|_{L^\infty} + \sqrt{\mathcal{H}(g(0))} < \delta$$

and $\|\nabla_{x,v}g(0)\|_{L^2} < +\infty$, then the initial value problem for (1.17) has a unique global in time solution satisfying the estimates (1.24) and (1.25).

Remark 1.3 This implies the stability of the mixed phase for $\beta < 1$. The proof is again a simpler version of the one for Theorem 1.2 which will be omitted. In this case the analog of operator A has a spectral gap property (because $\beta < 1$) and a trivial null space, so that we do not need to require the symmetry property (1.19).

We have already noticed that, when $\beta > 1$, the couple $(1, 1)$ is a local maximizer for the free energy. In next theorem we establish the dynamical instability of M_{hom} .

Theorem 1.4 (Instability) *Assume $\beta > 1$. There exist constants $k_0 > 0$, $\theta > 0$, $C > 0$, $c > 0$ and a family of initial $\frac{2\pi}{k_0}$ -periodic data $f_i^\delta(0) = \mu_\beta + \sqrt{\mu_\beta} g_i^\delta(0) \geq 0$, with $g^\delta(0)$ satisfying (1.19) and*

$$\|\nabla_{x,v} g^\delta(0)\|_{L^2} + \|w g^\delta(0)\|_{L^\infty} \leq C\delta,$$

for δ sufficiently small, but the solution $g^\delta(t)$ to (1.1) satisfies

$$\sup_{0 \leq t \leq T^\delta} \|w g^\delta(t)\|_{L^\infty} \geq c \sup_{0 \leq t \leq T^\delta} \|g^\delta(t)\|_{L^2} \geq c\theta > 0.$$

Here the escape time is

$$T^\delta = \frac{1}{\operatorname{Re}\lambda} \ln \frac{\theta}{\delta}, \quad (1.26)$$

λ is the eigenvalue with the largest real part for the linearized Vlasov-Boltzmann system constructed in Theorem 5.1 with $\operatorname{Re}\lambda > 0$.

Remark 1.4 Note that $T^\delta \rightarrow \infty$ as $\delta \rightarrow 0$. We also observe that the critical value $\beta = 1$ escapes our analysis, since it is based on some strict inequalities which collapse for the critical β . Furthermore, the growing mode which we construct satisfies the symmetry condition (1.19). Hence the instability is not a consequence of the absence of such a symmetry.

To prove instability of the homogeneous state, we encounter a second major difficulty. Even though such a homogeneous equilibrium is not a local minimizer for the free-energy (1.11), it is not clear at all if such a property leads to a dynamical instability along the time evolution. As a matter of fact, the presence of (possibly strong) stabilizing collision in the velocity space might damp out the instability. It is thus fundamental to understand such a stabilizing collision effect. Unfortunately, a direct linear stability analysis in the presence of the collision effect is too complicated to draw any conclusion. We therefore developed a new perturbation argument to establish the linear instability around such a homogeneous steady state. We first observe that in the absence of the collision effect, the homogeneous state is indeed dynamically unstable, by an explicit analysis similar to the Penrose criterion in plasma physics [14]. It then follows, via a perturbation argument, that for ‘weak’ collision effects, such an instability should persist. The key is to show that this is even true for arbitrarily ‘strong’ collision effect. We use an argument of contradiction and a method of continuation. In fact, if instability should fail at some level of the collision effect, then a neutral mode must occur. The interaction between the Vlasov force and the collision effect forces any neutral mode to behave like a multiple of M_{hom} with a particular dispersion relation never satisfied for our model (Theorem 5.1). To bootstrap such a linear instability into a nonlinear one is delicate due to severe nonlinearity, and we follow the program developed by Strauss and the second author over the years [10, 11].

We remark that the instability of the homogeneous equilibrium for VFP model is still open because the techniques used in the present paper cannot apply due to the unboundedness of the Fokker-Plank operator. On the other

hand, in this paper we cannot prove the convergence to the stable equilibrium, while it was possible in the VFP model even to compute the rate.

The paper is organized as follows: In Section 2 we use the energy-entropy to derive a mixed $L^1 - L^2$ estimate. In Section 3, we establish some lemmas on the characteristics curves for the equations (1.1). In Section 4 we establish the nonlinear stability in L^∞ norm. In Section 5 we construct a growing mode for the linearized problem around the homogeneous equilibrium and finally, in Section 6 we show that such a linear growing mode leads to nonlinear instability.

2 Entropy-Energy Estimate.

In this section we use the conservation of energy and masses, and the entropy inequality to obtain a priori estimates on the deviation of the solution from equilibrium. A crucial role is played by the quadratic approximation of the excess free energy functional $\hat{\mathcal{F}}$. We need a few definitions and some notation.

Let $u = (u_1, u_2)$ be a couple of functions in $L^2(\mathbb{R})$ and denote $\langle u, v \rangle = \sum_{i=1}^2 \langle u_i, v_i \rangle$. Given a couple of densities $\rho = (\rho_1(x), \rho_2(x))$, we define the operator A by setting

$$\langle u, Au \rangle := \sum_{i=1}^2 \int_{\mathbb{R}} dx u_i(x) (Au)_i(x) = \frac{1}{2} \frac{d^2}{ds^2} \hat{\mathcal{F}}(\rho + su)|_{s=0}.$$

In particular, when $\rho = (\bar{\rho}_1(x), \bar{\rho}_2(x))$, the action of the operator A on u is

$$(Au)_1 = \frac{u_1}{\bar{\rho}_1} + \beta U * u_2, \quad (Au)_2 = \frac{u_2}{\bar{\rho}_2} + \beta U * u_1. \quad (2.1)$$

Hence

$$\langle u, Au \rangle = \frac{1}{2} \sum_{i=1}^2 \int_{\mathbb{R}} dx \frac{u_i^2(x)}{\bar{\rho}_i(x)} + \beta \int_{\mathbb{R}} dx \int_{\mathbb{R}} dy U(|x-y|) u_1(x) u_2(y).$$

Due to the minimizing properties of $\bar{\rho}$, this quadratic form is non negative. Moreover, by (1.12) and (2.1), we see that $A\bar{\rho}' = 0$, which shows that $\bar{\rho}'$ is in the null space of A . Indeed, one can show (see [5] and references quoted therein) the following

Lemma 2.1 *Suppose $\beta > 1$ and $\rho = (\bar{\rho}_1, \bar{\rho}_2)$. Then there exist $\delta_0 > 0$ such that*

$$\langle u, Au \rangle \geq \delta_0 \langle (I - \mathcal{P})u, (I - \mathcal{P})u \rangle,$$

where \mathcal{P} is the projector on Null A :

$$\text{Null } A = \{u \in L^2(\mathbb{R}) \times L^2(\mathbb{R}) \mid u = c\bar{\rho}', c \in \mathbb{R}\}.$$

If either $\rho = (\rho^\pm, \rho^\mp)$, or $\rho = (1, 1)$ with $\beta < 1$, we have

$$\langle u, Au \rangle \geq \delta_0 \langle u, u \rangle,$$

and the null space of A reduces to $\{0\}$.

We prove the following lemma which plays a crucial role in the proof of the stability of $M_{\bar{\rho}}$. Remind that we have adopted the notation $\xi = (v, \zeta)$, with $\zeta \in \mathbb{R}^2$ and $v \in \mathbb{R}$ for the velocity.

Lemma 2.2 *Let $M = M_{\bar{\rho}}$. Let $f_1(t, x, v, \zeta) = f_2(t, -x, -v, \zeta)$ and that $\|g\|_{L^\infty} \leq \delta$ for some small δ . Then there exists $C > 0$ and $\kappa > 0$, such that*

$$C \sum_{i=1,2} \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi \left\{ \frac{(f_i(t) - M_i)^2}{M_i} \mathbf{1}_{|f_i(t) - M_i| \leq \kappa M_i} + f_i(t) - M_i \mathbf{1}_{|f_i(t) - M_i| \geq \kappa M_i} \right\} \leq \mathcal{H}(g(0)).$$

Remark 2.1 Note that, when dealing with M_{red} , M_{blue} and M_{hom} , the functions $\bar{\rho}_1$ and $\bar{\rho}_2$ have to be replaced by the constant values (ρ^+, ρ^-) , (ρ^-, ρ^+) and $(1, 1)$ respectively. With this modifications the lemma still holds.

Proof Remember the notation $\rho_f(t, x) = \int_{\mathbb{R}^3} d\xi f(t, x, \xi)$. We may construct solutions (see [8]) such that

$$\mathcal{H}(g) \leq \mathcal{H}(g(0)).$$

We expand $\mathcal{H}(g)$ and use (1.7) to cancel the linear part of the expansion, which takes the form

$$\begin{aligned} & - \sum_{i=1}^2 \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi \left\{ C_i + 1 + \ln\left(\frac{\beta}{2\pi}\right)^{\frac{3}{2}} \right\} \{f_i - M_i\} + \\ & \sum_{i=1}^2 \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi \left\{ \frac{\beta|\xi|^2}{2} (f_i - M_i) + \{\ln M_i + 1\} (f_i - M_i) \right. \\ & \left. + \beta \{U * \bar{\rho}_{i+1}\} (f_i - M_i) \right\}. \end{aligned}$$

Indeed, since $\ln M_i = \ln\left(\frac{\beta}{2\pi}\right)^{\frac{3}{2}} - \frac{\beta|\xi|^2}{2} + \ln \bar{\rho}_i$, by (1.7), the above quantity is zero by construction. Therefore, we turn to the second order expansion of $\mathcal{H}(g)$. For some \tilde{f}_i between M_i and f_i ,

$$\begin{aligned} \mathcal{H}(g) &= \sum_{i=1}^2 \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi \frac{(f_i(t) - M_i)^2}{2\tilde{f}_i} dx d\xi \\ &+ \beta \int_{\mathbb{R}} dx \int_{\mathbb{R}} dy (\rho_{f_1}(t, x) - \bar{\rho}_1(x)) U(|x - y|) (\rho_{f_2}(t, x) - \bar{\rho}_2(x)). \end{aligned} \quad (2.2)$$

For some small number κ to be determined, we introduce the indicator functions $\chi_i^< = \mathbf{1}_{|f_i(t) - M_i| \leq \kappa M_i}$ and $\chi_i^> = \mathbf{1}_{|f_i(t) - M_i| > \kappa M_i}$ and split the first term into

$$\int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi \frac{(f_i(t) - M_i)^2}{2\tilde{f}_i} \chi_i^> + \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi \frac{(f_i(t) - M_i)^2}{2\tilde{f}_i} \chi_i^<$$

We first estimate $\frac{(f_i(t) - M_i)^2}{2\tilde{f}_i}$ in the case of $|f_i(t) - M_i| > \kappa M_i$. Notice that either $f_i \geq (1 + \kappa)M_i$, or $f_i \leq (1 - \kappa)M_i$. If $f_i \geq (1 + \kappa)M_i$, $\tilde{f}_i(t) \leq f_i$, and we have

$$\frac{|f_i(t) - M_i|}{\tilde{f}_i(t)} \geq \frac{|f_i(t) - M_i|}{f_i} = 1 - \frac{M_i}{f_i} \geq 1 - \frac{1}{1 + \kappa} = \frac{\kappa}{1 + \kappa}.$$

In the second case $f_i \leq (1 - \kappa)M_i$, $\tilde{f}_i(t) \leq M_i$ and

$$\frac{|f_i(t) - M_i|}{\tilde{f}_i(t)} \geq \frac{|f_i(t) - M_i|}{M_i} = 1 - \frac{f_i}{M_i} \geq 1 - \{1 - \kappa\} = \kappa > \frac{\kappa}{1 + \kappa}.$$

Combining these two cases and noticing $\tilde{f}_i \leq (1 + \kappa)M_i$ for $|f_i(t) - M_i| \leq \kappa M_i$, we conclude

$$\begin{aligned} & \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi \frac{(f_i(t) - M_i)^2}{2\tilde{f}_i} \chi_i^{\leq} + \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi \frac{(f_i(t) - M_i)^2}{2\tilde{f}_i} \chi_i^{\geq} \\ & \geq \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi \frac{(f_i(t) - M_i)^2}{2(1 + \kappa)M_i} \chi_i^{\leq} + \frac{\kappa}{2(1 + \kappa)} \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi |f_i(t) - M_i| \chi_i^{\geq} \\ & = \frac{1}{2(1 + \kappa)} \int \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi g_i^2 \chi_i^{\leq} + \frac{\kappa}{2(1 + \kappa)} \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi |f_i(t) - M_i| \chi_i^{\geq} \\ & \geq \frac{1}{2(1 + \kappa)} \int \int_{\mathbb{R}} dx n_i^2 + \frac{\kappa}{2(1 + \kappa)} \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi |f_i(t) - M_i| \chi_i^{\geq}, \end{aligned} \tag{2.3}$$

where we have set $n_i(t, x) \sqrt{\mu_\beta} = \mathbf{P}[(f_i(t) - M_i) \chi_i^{\leq}]$, and \mathbf{P} denotes the L_ξ^2 -projection on $\sqrt{M_i}$:

$$\mathbf{P}f := \left(\int_{\mathbb{R}^3} d\xi' f(\xi') \sqrt{\mu_\beta(\xi')} \right) \sqrt{\mu_\beta(\xi)}.$$

We now split the potential contribution in (2.3) to get

$$\begin{aligned} & \beta \int_{\mathbb{R}} dx \int_{\mathbb{R}} dy \int_{\mathbb{R}^3} d\xi \int_{\mathbb{R}^3} d\eta \sqrt{M_1} g_1 \chi_1^{\leq} U(|x - y|) \sqrt{M_2} g_2 \chi_2^{\leq} \\ & + \beta \int_{\mathbb{R}} dx \int_{\mathbb{R}} dy \int_{\mathbb{R}^3} d\xi \int_{\mathbb{R}^3} d\eta \sqrt{M_1} g_1 \chi_1^{\leq} U(|x - y|) \sqrt{M_2} g_2 \chi_2^{\geq} \\ & + \beta \int_{\mathbb{R}} dx \int_{\mathbb{R}} dy \int_{\mathbb{R}^3} d\xi \int_{\mathbb{R}^3} d\eta \sqrt{M_1} g_1 \chi_1^{\geq} U(|x - y|) \sqrt{M_2} g_2 \chi_2^{\geq} \\ & + \beta \int_{\mathbb{R}} dx \int_{\mathbb{R}} dy \int_{\mathbb{R}^3} d\xi \int_{\mathbb{R}^3} d\eta \sqrt{M_1} g_1 \chi_1^{\geq} U(|x - y|) \sqrt{M_2} g_2 \chi_2^{\leq}. \end{aligned}$$

From our assumption $\|g\|_{L^\infty} \leq \delta$, the last three terms are controlled by

$$\begin{aligned} & C_\beta \left(\|g_1\|_{L^\infty} + \|g_2\|_{L^\infty} \right) \sum_{i=2} \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi |f_i(t) - M_i| \chi_i^{\geq} \\ & \leq C_\beta \delta \sum_{i=1}^2 \int_{\mathbb{R}} dx \int_{\mathbb{R}^3} d\xi |f_i(t) - M_i| \mathbf{1}_{|f_i(t) - M_i| \geq \kappa M_i}. \end{aligned}$$

which is bounded by the second term in (2.3) for $\delta \ll \kappa$. Since $\int d\xi \sqrt{M_i} g_i = \sqrt{\bar{\rho}_i} n_i$, the first term can be written as

$$\beta \int_{\mathbb{R}} dx \int_{\mathbb{R}} dy n_1(t, x) \sqrt{\bar{\rho}_1(x)} U(|x - y|) n_2(t, y) \sqrt{\bar{\rho}_2(y)}$$

We now combine it with the first term in (2.3) to get

$$\begin{aligned} & \frac{1}{2(1 + \kappa)} \int_{\mathbb{R}} dx (n_1^2 + n_2^2) \\ & + \beta \int_{\mathbb{R}} dx \int_{\mathbb{R}} dy n_1(x) \sqrt{\bar{\rho}_1(x)} U(|x - y|) n_2(y) \sqrt{\bar{\rho}_2(y)} dx dy \\ & = \langle n \sqrt{\bar{\rho}}, An \sqrt{\bar{\rho}} \rangle - \frac{\kappa}{2(1 + \kappa)} \int_{\mathbb{R}} dx (n_1^2 + n_2^2). \end{aligned}$$

Since $n_1(x) = n_2(-x)$ by our symmetry assumption and from $\bar{\rho}'_1(x) = -\bar{\rho}'_2(-x)$, it follows that (n_1, n_2) is orthogonal to the null space of A . From the spectral inequality for the operator A , since $\bar{\rho}_i(x) \geq \rho^-$, the above is bounded from below by

$$\begin{aligned} & \left(\rho^- \delta_0 - \frac{\kappa}{2(1 + \kappa)} \right) \int_{\mathbb{R}} dx (n_1^2 + n_2^2) \\ & \geq c \sum_{i=1}^2 \|\mathbf{P}(f_i(t) - M_i) \mathbf{1}_{|f_i(t) - M_i| \leq \kappa M_i}\|_{L^2}^2, \end{aligned}$$

provided that we choose κ sufficiently small. Then the lemma follows by collecting the terms with $\delta \ll \kappa$.

3 1-Dimensional Characteristics.

We define the characteristics curves $[X_i(s; t, x, v), V_i(s; t, x, v)]$ for (1.17) passing through (t, x, v) at $s = t$, such that

$$\begin{aligned} \frac{dX_i(s; t, x, v)}{ds} &= V_i(s; t, x, v), \\ \frac{dV_i(s; t, x, v)}{ds} &= -\partial_x U * \int_{\mathbb{R}^3} d\xi (M_{i+1} + \sqrt{M_{i+1} g_{i+1}}) \equiv -\partial_x \phi(X_i(s; t, x, v)). \end{aligned} \quad (3.1)$$

We also define the unperturbed characteristics $[X_i^0(s; t, x, v), V_i^0(s; t, x, v)]$ passing through (t, x, v) at $s = t$, such that

$$\begin{aligned} \frac{dX_i^0(s; t, x, v)}{ds} &= V_i^0(s; t, x, v), \\ \frac{dV_i^0(s; t, x, v)}{ds} &= -\partial_x U * \int M_{i+1} d\xi \equiv -\partial_x \phi_0(X_i^0(s; t, x, v)) \end{aligned} \quad (3.2)$$

Our main goal is to study the zero set of $\frac{\partial X_i(s; t, x, v)}{\partial v}$.

Lemma 3.1 For any (t, x, v) with $v \neq 0$, the set of $\{s \in \mathbb{R} : \frac{\partial X_i^0(s; t, x, v)}{\partial v} = 0\}$ is countable.

Proof From the particle energy conservation for (3.2):

$$\frac{1}{2}|V_i^0(s; t, x, v)|^2 - \phi_0(X_i^0(s; t, x, v)) = \frac{1}{2}|v|^2 - \phi_0(x). \quad (3.3)$$

Taking derivative with respect to v yields

$$V_i^0(s; t, x, v) \frac{\partial V_i^0(s; t, x, v)}{\partial v} - \phi_0' \frac{\partial X_i^0(s; t, x, v)}{\partial v} = v.$$

Assume $v \neq 0$. If there is s_0 such that $\frac{\partial X_i^0(s_0; t, x, v)}{\partial v} = 0$, then necessarily

$$V_i^0(s_0; t, x, v) \frac{\partial V_i^0(s_0; t, x, v)}{\partial v} \neq 0,$$

because, for such an s_0 it reduces to v by previous equation and we assumed $v \neq 0$. Hence

$$\frac{\partial V_i^0(s_0; t, x, v)}{\partial v} = \frac{d}{ds} \frac{\partial X_i^0(s; t, x, v)}{\partial v} \Big|_{s=s_0} \neq 0$$

and therefore $\frac{\partial X_i^0(s; t, x, v)}{\partial v} \neq 0$ in a neighborhood of $s = s_0$. This implies that for $v \neq 0$, the set $\{s : \frac{\partial X_i^0(s; t, x, v)}{\partial v} = 0\}$ is countable.

Lemma 3.2 Fix $T_0 > 0$ and $N > 0$. Let $|v| \leq N$.

1. For any $\varepsilon > 0$, there exists L_ε sufficiently large so that, if $|x| \geq L_\varepsilon$, then for $0 \leq s \leq t - \varepsilon$,

$$\frac{\partial X_i^0(s; t, x, v)}{\partial v} < -\frac{\varepsilon}{2} < 0.$$

2. For any $\eta > 0$, there exist P finite points $|x_k| \leq L_\varepsilon$ ($1 \leq k \leq P$) and corresponding open sets

$$O_{x_k} = \bigcup_{(n,o) \in I_k} \{a_n < s < b_n\} \times \{c_o < v < d_o\}$$

with the property

$$|[0, T_0] \times \{|v| \leq N\} \cap O_{x_k}^c| < \eta,$$

so that there exists $m > 0$ and, for any $|x| \leq L_\varepsilon$, there exists $l \in \{1, \dots, P\}$

$$\left| \frac{\partial X_i^0(s; t, x, v)}{\partial v} \right| > m > 0.$$

for $(s, v) \in O_{x_l}$.

Proof For any $\varepsilon > 0$, from (3.2) and (3.3) ,

$$\begin{aligned} |V_i^0(s; t, x, v)| &\leq |v| + 2\sqrt{\|\phi_0\|_{L^\infty}} \leq N + C, \\ |X_i^0(s; t, x, v) - x| &\leq T_0\{N + C\}. \end{aligned}$$

By choosing L_ε (depending on T_0 and N) large enough, for $|x| \geq L_\varepsilon$, $|v| \leq N$, and $0 \leq s \leq T_0$,

$$|X_i^0(s; t, x, v)| \geq \frac{L_\varepsilon}{2}. \quad (3.4)$$

From (3.2), we have

$$\frac{d^2}{ds^2} \frac{\partial X_i^0(s; t, x, v)}{\partial v} = -\partial_{xx}\phi_0(X_i^0(s; t, x, v)) \frac{\partial X_i^0(s; t, x, v)}{\partial v} \quad (3.5)$$

and we deduce that for $|s| \leq T_0$,

$$\left| \frac{\partial X_i^0(s; t, x, v)}{\partial v} \right| \leq C_{T_0}. \quad (3.6)$$

By the Taylor expansion for s , we get

$$\begin{aligned} \frac{\partial X_i^0(s; t, x, v)}{\partial v} &= \frac{\partial X_i^0(s; t, x, v)}{\partial v} \Big|_{s=t} + (s-t) \frac{d}{ds} \frac{\partial X_i^0(s; t, x, v)}{\partial v} \Big|_{s=t} \\ &\quad + \frac{(s-t)^2}{2} \frac{d^2}{ds^2} \frac{\partial X_i^0(\bar{s}; t, x, v)}{\partial v} \\ &= (s-t) + \frac{(s-t)^2}{2} \frac{d^2}{ds^2} \frac{\partial X_i^0(\bar{s}; t, x, v)}{\partial v} \end{aligned}$$

for some $t - T_0 \leq \bar{s} \leq t$. Since the densities $\bar{\rho}_i$ tend to their asymptotic values at infinity, $\lim_{y \rightarrow \infty} \partial_{xx}\phi_0(y) = 0$. Therefore, using again (3.5) with $s = \bar{s}$ and (3.6), we have

$$\begin{aligned} \frac{\partial X_i^0(s; t, x, v)}{\partial v} &\leq (s-t) \left(1 - (t-s)C_{T_0} \sup_{|y| \geq \frac{L_\varepsilon}{2}} |\partial_{xx}\phi_0(y)|\right) \\ &\leq \frac{s-t}{2} < 0, \end{aligned}$$

by choosing L_ε sufficiently large. Part (1) thus follows.

To prove part (2), for $|x| \leq L_\varepsilon$, introduce the zero set of $\frac{\partial X_i^0(s; t, x, v)}{\partial v}$ as

$$Z_x = \{t - T_0 \leq s \leq T_0, |v| \leq N : \frac{\partial X_i^0(s; t, x, v)}{\partial v} = 0\}.$$

Then from the Fubini Theorem and Lemma 3.1,

$$|Z_x| = \int_{-N}^N \left\{ \int_{t-T_0}^t \mathbf{1}_{\{s, v: \frac{\partial X_i^0(s; t, x, v)}{\partial v} = 0\}} ds \right\} dv = 0.$$

Therefore, there exists an open set Ω_x such that $Z_x \subset \Omega_x$ with $|\Omega_x| < \frac{\eta}{2}$.

Clearly $\frac{\partial X_i^0(s; t, x, v)}{\partial v} \neq 0$ over the compact set $[t - T_0, t] \times \{|v| \leq N\} \cap \Omega_x^c$. By

the continuity in s and v , there exists $m_x > 0$ such that over $[t - T_0, t] \times \{|v| \leq N\} \cap \Omega_x^c$

$$\left| \frac{\partial X_i^0(s; t, x, v)}{\partial v} \right| > 4m_x > 0.$$

Furthermore, from the continuity in x , we have an open set of $(x - \Delta_x, x + \Delta_x)$ such that

$$\left| \frac{\partial X_i^0(s; t, x', v)}{\partial v} \right| > 2m_x > 0$$

for all $x' \in (x - \Delta_x, x + \Delta_x)$, $(s, v) \in [t - T_0, t] \times \{|v| \leq N\} \cap \Omega_x^c$. Such $(x - \Delta_x, x + \Delta_x)$ forms an open covering for $|x| \leq L_\varepsilon$, hence there is a finite subcovering $\{(x_k - \Delta_k, x_k + \Delta_k), k = 1, \dots, P\}$ for $|x| \leq L_\varepsilon$. For any $|x| \leq L_\varepsilon$, there exists Ω_{x_l} such that $x \in (x_l - \Delta_l, x_l + \Delta_l)$ and for $[t - T_0, t] \times \{|v| \leq N\} \cap \Omega_{x_l}^c$,

$$\left| \frac{\partial X_i^0(s; t, x, v)}{\partial v} \right| > 2m = \min_{1 \leq k \leq P} 2m_k > 0.$$

We finally choose an (finite) open covering of $[t - T_0, t] \times \{|v| \leq N\} \cap \Omega_{x_k}^c$ of the form $O_{x_k} = \bigcup_{n,o} \{a_n < s < b_n\} \times \{c_o < v < d_o\}$ with $|a_n - b_n| + |c_o - d_o|$ sufficiently small so that over O_{x_k}

$$\left| \frac{\partial X_i^0(s; t, x, v)}{\partial v} \right| > m > 0$$

Lemma 3.3 Fix $T_0 > 0$ and $N > 0$. Let $|v| \leq N$. For any $\varepsilon > 0$, recall L_ε , O_{x_k} , $1 \leq k \leq P$ constructed in Lemma 3.2. There exists $\delta > 0$ such that if $\|g\|_{L^2} < \delta$,

1. If $|x| \geq L_\varepsilon$ then for $0 \leq s \leq t - \varepsilon$,

$$\frac{\partial X_i(s; t, x, v)}{\partial v} < -\frac{\varepsilon}{2}.$$

2. For any $|x| \leq L_\varepsilon$, there exists l such that $x \in (x_l - \Delta_l, x_l + \Delta_l)$ and for all $(s, v) \in O_{x_l}$

$$\left| \frac{\partial X_i(s; t, x, v)}{\partial v} \right| > \frac{m}{2} > 0.$$

Proof Denote the solution operator of (3.2) and (3.1) by G_0 and G . By the Duhamel principle, we have

$$\begin{aligned} \begin{pmatrix} X_i(s; t, x, v) \\ V_i(s; t, x, v) \end{pmatrix} &= e^{G_0(s-t)} \begin{pmatrix} x \\ v \end{pmatrix} \\ &+ \int_t^s e^{G_0(s-\tau)} \begin{pmatrix} 0 \\ -\partial_x U * \left\{ \int \sqrt{M_{i+1}} g_{i+1} d\xi^l \right\}(\tau) \end{pmatrix} d\tau. \end{aligned} \quad (3.7)$$

Notice that

$$\begin{aligned} & \left| \frac{\partial}{\partial v} \partial_x U * \left\{ \int \sqrt{M_{i+1} g_{i+1}} d\xi' \right\}(\tau) \right| \\ &= \left| \partial_{xx} U * \left\{ \int \sqrt{M_{i+1} g_{i+1}} d\xi' \right\} \frac{\partial X_i(\tau; t, x, v)}{\partial v} \right| \\ &\leq C \|g\|_{L^2} \left| \frac{\partial X_i(\tau; t, x, v)}{\partial v} \right|. \end{aligned}$$

It thus follows from taking derivative of v in (3.7) and by the Gronwall's lemma that $0 \leq s \leq T_0$

$$\left| \frac{\partial X_i(s; t, x, v)}{\partial v} \right| \leq e^{CT_0}.$$

We now use (3.7) again to get

$$\left| \frac{\partial X_i(s; t, x, v)}{\partial v} - \frac{\partial X_i^0(s; t, x, v)}{\partial v} \right| \leq C_{T_0} \|g\|_{L^2}.$$

Hence, we deduce our lemma by choosing $\|g\|_{L^2}$ sufficiently small.

4 Weighted L^∞ Stability.

In this section we use the entropy-energy bound and the estimates on the characteristics to show that the perturbation g of the non homogeneous equilibrium $M_{\bar{\rho}}$, is arbitrarily small at any positive time in a suitable weighted L^∞ norm, provided that it is initially sufficiently small, thus showing the stability of the non homogeneous equilibrium. We use the weight function $w(\xi) = (\Sigma + |\xi|^2)^\gamma$, with Σ a positive constant to be chosen later and $\gamma > \frac{3}{2}$.

Lemma 4.1 *Let $h = wg$. There exist $T_0 > 0$ and $\delta > 0$ such that, if $\|h\|_{L^\infty} < \delta$, then*

$$\|h(T_0)\|_{L^\infty} \leq \frac{1}{2} \|h(0)\|_{L^\infty} + C_{T_0} \sqrt{\mathcal{H}(g(0))}.$$

Proof We first write the equation for $h = wg$ from (1.17):

$$\begin{aligned} & \left(\partial_t + v \partial_x + F(M_{i+1} + \sqrt{M_{i+1}} \frac{h_{i+1}}{w}) \right) \partial_v + \nu(x, \xi) h_i = \\ & F \left(M_{i+1} + \sqrt{M_{i+1}} \frac{h_{i+1}}{w} \right) \frac{w'}{w} h_i + w F \left(\sqrt{M_{i+1}} g_{i+1} \right) v \sqrt{M_i} + \sum_{j=1,2} K_w^{ij} h_j \\ & + F \left(\sqrt{M_{i+1}} \frac{h_{i+1}}{w} \right) v h_i + w \Gamma \left(\frac{h_i}{w}, \frac{h_i}{w} \right) + w \Gamma \left(\frac{h_i}{w}, \frac{h_{i+1}}{w} \right), \end{aligned} \quad (4.1)$$

where $K_w^{ij}(\cdot) = w K^{ij} \left(\frac{\cdot}{w} \right)$. We note that for $\Sigma \geq 1$

$$\frac{w(\xi)}{w(\xi')} = \frac{[\Sigma + |\xi|^2]^\gamma}{[\Sigma + |\xi'|^2]^\gamma} \leq C_\gamma \frac{[\Sigma + |\xi'|^2]^\gamma + |\xi' - \xi|^{2\gamma}}{[\Sigma + |\xi'|^2]^\gamma} \leq C_\gamma [1 + |\xi' - \xi|^2]^\gamma. \quad (4.2)$$

For any (t, x, ξ) , integrating along its backward trajectory (3.1)

$$[X_i(s), V_i(s)] = [X_i(s; t, x, v), V_i(s; t, x, v)]$$

we can express $h_i(t, x, \xi)$ as

$$\begin{aligned} & h_i(0, X_i(0; t, x, v), V_i(0; t, x, v), \zeta) + \\ & \int_0^t e^{\int_t^s \nu_i(\tau) d\tau} \{F(M_{i+1} + \sqrt{M_{i+1}g_{i+1}}) \frac{w'}{w} h_i\}(s, X_i(s), V_i(s), \zeta) ds \\ & + \int_0^t e^{\int_t^s \nu_i(\tau) d\tau} \{F(\sqrt{M_{i+1}g_{i+1}}) V_i(s) \sqrt{M_i}\}(s, X_i(s), V_i(s), \zeta) ds \\ & + \sum_{j=1}^2 \int_0^t e^{\int_t^s \nu_i(\tau) d\tau} \left(\sum_{j=1,2} K_w^{i,j} h_j \right) (s, X_i(s), V_i(s), \zeta) ds \\ & + \int_0^t e^{\int_t^s \nu_i(\tau) d\tau} \{F(\sqrt{M_{i+1}g_{i+1}}) v h_i\}(s, X_i(s), V_i(s), \zeta) ds \\ & + \int_0^t e^{\int_t^s \nu_i(\tau) d\tau} w \left[\Gamma\left(\frac{h_i}{w}, \frac{h_i}{w}\right) + \Gamma\left(\frac{h_i}{w}, \frac{h_{i+1}}{w}\right) \right] (s, X_i(s), V_i(s), \zeta) ds \end{aligned} \quad (4.3)$$

We have set $\nu_i(\tau) \equiv \nu(V_i(\tau), \zeta) \geq \nu_0 > 0$. Fix a small constant $\varepsilon > 0$. We can choose Σ large so that $|\frac{w'}{w}| \leq \varepsilon$. Since $\|F(M_{i+1} + \sqrt{M_{i+1}g_{i+1}})\|_{L^\infty} \leq C$ if $\|h\|_{L^\infty}$ is small, the second term in (4.3) is bounded by

$$C\varepsilon e^{-\frac{\nu_0 t}{2}} \sup_{0 \leq s \leq T_0} \{e^{\frac{\nu_0 s}{2}} \|h(s)\|_{L^\infty}\}.$$

For the third term in (4.3), we split

$$g_{i+1} = g_{i+1} \mathbf{1}_{|f_i(t) - M_i| \geq \kappa M_i} + g_{i+1} \mathbf{1}_{|f_i(t) - M_i| < \kappa M_i}.$$

Since U is smooth,

$$\begin{aligned} & \|F(\sqrt{M_{i+1}g_{i+1}})\|_{L^\infty} \\ & \leq C \{ \|\sqrt{M_{i+1}g_{i+1}} \mathbf{1}_{|f_i(t) - M_i| \leq \kappa M_i}\|_{L^2} + \|\sqrt{M_{i+1}g_{i+1}} \mathbf{1}_{|f_i(t) - M_i| \geq \kappa M_i}\|_{L^1} \}, \end{aligned}$$

by Lemma 2.2,

$$\begin{aligned} & \|F(\sqrt{M_{i+1}g_{i+1}}) V_i(s) \sqrt{M_i}(s, X_i(s), V_i(s), \zeta)\|_{L^\infty} \\ & \leq C(\mathcal{H}(g(0)) + \sqrt{\mathcal{H}(g(0))}). \end{aligned} \quad (4.4)$$

For the fifth term, we note that, since for hard spheres $\nu_i(s) \geq \nu_0(1 + |\zeta| + |V_i(s)(s)|)$, it follows that $\frac{|V_i(s)|}{\nu_i(s)} < \frac{1}{\nu_0}$. Moreover, $\int_0^t ds \frac{\nu_i(s)}{2} e^{\int_t^s (\nu_i(\tau)/2) d\tau} \leq 1$. Therefore,

$$\left| \int_0^t e^{\int_t^s \nu_i(\tau) d\tau} F(\sqrt{M_{i+1}g_{i+1}}) V_i(s) h_i ds \right| \leq C e^{-\nu_0 t} \sup_{0 \leq s \leq T_0} \{e^{\frac{\nu_0 s}{2}} \|h(s)\|_{L^\infty}\}^2.$$

For the last term in (4.3), by Lemma 10 of [7], it follows

$$\left| w\Gamma\left(\frac{h_i}{w}, \frac{h_i}{w}\right)(\xi) \right| + \left| w\Gamma\left(\frac{h_i}{w}, \frac{h_{i+1}}{w}\right)(\xi) \right| \leq C\nu(\xi)\|h\|_{L^\infty}^2.$$

We therefore get the bound for the last term by

$$\begin{aligned} & \int_0^t e^{\int_t^s \nu(\tau) d\tau} \nu_i(s) \|h(s)\|_{L^\infty}^2 ds \\ & \leq C \left\{ \sup_{0 \leq s \leq T_0} e^{\frac{\nu_0 s}{2}} \|h(s)\|_{L^\infty} \right\}^2 \int_0^t e^{\int_t^s \nu_i(\tau) d\tau} \nu_i(s) e^{-\nu_0 s} ds. \end{aligned}$$

Note that $\frac{d}{ds}[e^{\int_t^s \nu_i(\tau) d\tau}] = e^{\int_t^s \nu_i(\tau) d\tau} \nu_i(s)$. Integrating by parts yields

$$\begin{aligned} \int_0^t e^{\int_t^s \nu_i(\tau) d\tau} \nu_i(s) e^{-\nu_0 s} ds &= \left(e^{\int_t^s \nu_i(\tau) d\tau} e^{-\nu_0 s} \right) \Big|_{s=0}^{s=t} + \nu_0 \int_0^t e^{\int_t^s \nu_i(\tau) d\tau} e^{-\nu_0 s} ds \\ &\leq C(1+t)e^{-\nu_0 t}. \end{aligned}$$

We shall mainly concentrate on the fourth term in (4.3). Let $\mathbf{k}^{i,j}(\xi, \xi')$ be the corresponding kernel associated with K_w^{ij} in (4.1). We now use (4.3) for $h_j(s, X_i(s), \xi')$ again to evaluate

$$\{K_w^{i,j} h_j\}(s, X_i(s), V_i(s), \zeta) = \int \mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi') h_j(s, X_i(s), \xi') d\xi'.$$

Denote

$$[X_j(s_1), V_j(s_1)] \equiv [X_j(s_1; X_i(s; t, x, v), v'), V_j(s_1; X_i(s; t, x, v), v')].$$

We can bound the fourth term in (4.3) by the sum on j of

$$\begin{aligned} & \int_0^t e^{\int_t^s \nu_i(\tau) d\tau + \int_s^0 \nu_j(\tau) d\tau} \int_{\mathbb{R}^3} |\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi')| h_j(0, X_j(0), V_j(0), \zeta') |d\xi' ds \\ & + \int_0^t \int_{s_1}^s e^{\int_t^s \nu_i(\tau) d\tau + \int_s^{s_1} \nu_j(\tau) d\tau} \int_{\mathbb{R}^3} |\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi')| \\ & \times \{F(M_{j+1} + \sqrt{M_{j+1} g_{j+1}}) \frac{w'}{w} h_j\}(s_1, X_j(s_1), V_j(s_1), \zeta') d\xi' ds ds_1 \\ & + \int_0^t \int_{s_1}^s e^{\int_t^s \nu_i(\tau) d\tau + \int_s^{s_1} \nu_j(\tau) d\tau} \int_{\mathbb{R}^3} |\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi')| \\ & \times \{F(\sqrt{M_{j+1} g_{j+1}}) v \sqrt{M_j}\}(s_1, X_j(s_1), V_j(s_1), \zeta') d\xi' ds \\ & + \sum_k \int_0^t \int_{s_1}^s \int_{\mathbb{R}^3 \times \mathbb{R}^3} e^{\int_t^s \nu_i(\tau) d\tau + \int_s^{s_1} \nu_j(\tau) d\tau} |\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi') \\ & \times \mathbf{k}_w^{j,k}(V_j(s_1), \zeta', \xi'') |h_k(s_1, X_j(s_1), \xi'')| d\xi' d\xi'' ds ds_1 \end{aligned} \quad (4.5)$$

$$\begin{aligned}
& + \int_0^t \int_{s_1}^s e^{\int_t^s \nu_i(\tau) d\tau + \int_{s_1}^s \nu_j(\tau) d\tau} \int_{\mathbb{R}^3} |\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi')| \\
& \times \{F(\sqrt{M_{j+1}g_{j+1}})vh_j\}(s_1, X_j(s_1), V_j(s_1), \zeta') d\xi' ds ds_1 \\
& + \int_0^t \int_{s_1}^s e^{\int_t^s \nu_i(\tau) d\tau + \int_{s_1}^s \nu_j(\tau) d\tau} \int_{\mathbb{R}^3} |\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi')| \\
& \times w \left\{ \Gamma\left(\frac{h_j}{w}, \frac{h_j}{w}\right) + \Gamma\left(\frac{h_j}{w}, \frac{h_{j+1}}{w}\right) \right\}(s, X_j(s_1), V_j(s_1)) d\xi' ds ds_1.
\end{aligned}$$

We will make an extended use of Lemma 7 of [7], which we report here for reader's convenience: For hard spheres, the usual Grad estimates imply:

$$|\mathbf{k}_{i,j}(\xi, \xi')| \leq C \{|\xi - \xi'| + |\xi - \xi'|^{-1}\} e^{-\frac{1}{8}|\xi - \xi'|^2 - \frac{1}{8} \frac{\|\xi\|^2 - \|\xi'\|^2}{|\xi - \xi'|^2}}. \quad (4.6)$$

Lemma 4.2 (Lemma 7 of [7]) *There are $\varepsilon > 0$ and $C > 0$ such that*

$$\int_{\mathbb{R}^3} d\xi' \frac{w(\xi)}{w(\xi')} \{|\xi - \xi'| + |\xi - \xi'|^{-1}\} e^{-\frac{1-\varepsilon}{8}|\xi - \xi'|^2 - \frac{1-\varepsilon}{8} \frac{\|\xi\|^2 - \|\xi'\|^2}{|\xi - \xi'|^2}} \leq \frac{C}{1 + |\xi|}. \quad (4.7)$$

By Lemma 4.2, we obtain the crucial estimate

$$\int_{\mathbb{R}^3} \mathbf{k}_w^{i,j}(|\xi, \xi'|) d\xi' < \frac{C}{1 + |\xi|} \quad (4.8)$$

uniformly in Σ . Since $\nu_i \geq \nu_0$, by taking L^∞ norm for h and (4.8), we bound the first term in (4.6) by $Cte^{-\nu_0 t} \|h_0\|_{L^\infty}$, and the second term by

$$\varepsilon C e^{-\frac{\nu_0}{2} t} \sup_{0 \leq s \leq T_0} \{e^{\frac{\nu_0}{2} s} \|h(s)\|_{L^\infty}\}.$$

By (4.8), the third term is bounded by $C\mathcal{H}(g(0))$ as in (4.4), and the last two nonlinear terms are bounded by

$$C\{1 + t\} e^{-\nu_0 t} \left\{ \sup_{0 \leq s \leq T_0} e^{\frac{\nu_0}{2} s} \|h(s)\|_{L^\infty} \right\}^2.$$

We now concentrate on the fourth term in (4.6), which will be estimated along the same lines of the proof of Theorem 20 in [7].

CASE 1: For $|\xi| \geq N_{T_0}$, we know that from (3.1)

$$|V_i(s; t, x, v) - v| \leq |s - t|C \leq CT_0,$$

By Lemma 4.2 and (4.2), for N_{T_0} large

$$\begin{aligned}
& \int \int |\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi') \mathbf{k}_w^{j,k}(V_j(s_1), \zeta', \xi'')| d\xi' d\xi'' \\
& \leq \frac{C}{1 + |V_i(s)| + |\zeta|} \leq \frac{C}{N - CT_0},
\end{aligned}$$

we therefore can find an upper bound for the fourth term in (4.6) by ($N \gg T_0$)

$$\begin{aligned} \frac{C}{N} \int_0^t e^{-\nu_0(t-s)} \times \int_0^s e^{-\nu_0(s-s_1)} \|h(s_1)\|_{L^\infty} ds_1 ds \\ \leq \frac{C e^{-\frac{\nu_0}{2}t}}{N} \sup_{0 \leq s \leq T_0} e^{\frac{\nu_0}{2}s} \|h(s)\|_{L^\infty}. \end{aligned}$$

CASE 2: For $|\xi| \leq N$, $|\xi'| \geq 2N$, or $|\xi'| \leq 2N$, $|\xi''| \geq 3N$. Notice that we have either $|\xi' - \xi| \geq N$ or $|\xi' - \xi''| \geq N$. This implies that

$$\begin{aligned} |v' - V_i(s; t, x, v)| &\geq |v' - v| - |v - V_i(s; t, x, v)| \\ &\geq |v' - v| - CT_0, \\ |v'' - V_j(s_1; X_i(s; t, x, v), v')| &\geq |v'' - v'| - |v' - V_j(s_1; X_i(s; t, x, v), v')| \\ &\geq |v'' - v'| - CT_0. \end{aligned}$$

Therefore, either one of the following are valid correspondingly for some $\sigma > 0$:

$$\begin{aligned} |\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi')| &\leq C_{T_0} e^{-\frac{\sigma}{8}N^2} |\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi')| e^{\frac{\sigma}{8}\{|V_i(s)-v'|^2+|\zeta-\zeta'|^2\}}, \\ |\mathbf{k}_w^{j,k}(V_j(s_1), \zeta', \xi'')| &\leq C_{T_0} e^{-\frac{\sigma}{8}N^2} |\mathbf{k}_w^{j,k}(V_j(s_1), \zeta', \xi'')| \times \\ &\quad e^{\frac{\sigma}{8}\{|V_j(s_1)-v''|^2+|\zeta''-\zeta'|^2\}}. \end{aligned}$$

From Lemma 4.2,

$$\begin{aligned} \int |\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi')| e^{\frac{\sigma}{8}\{|V_i(s)-v'|^2+|\zeta-\zeta'|^2\}} |d\xi' \\ + \int |\mathbf{k}_w^{j,k}(V_j(s_1), \zeta', \xi'')| e^{\frac{\sigma}{8}\{|V_j(s_1)-v''|^2+|\zeta''-\zeta'|^2\}} |d\xi'' < +\infty. \end{aligned} \quad (4.9)$$

We use this bound to combine the cases of $|\xi' - \xi| \geq N$ or $|\xi' - \xi''| \geq N$ as:

$$\int_0^t \int_0^{s_1} \left\{ \int_{|\xi| \leq N, |\xi'| \geq 2N} + \int_{|\xi'| \leq 2N, |\xi''| \geq 3N} \right\}$$

We first integrate ξ' for the first integral and apply (4.8) to integrate $\mathbf{k}_w^{j,k}$ over ξ'' . We then integrate ξ'' for the second integral and apply (4.8) to integrate $\mathbf{k}_w^{i,j}$ over ξ' . We thus find an upper bound

$$\begin{aligned} C \int_0^t \int_0^{s_1} \sup_{\xi} \int_{|\xi| \leq N, |\xi'| \geq 2N} |\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi')| d\xi' \\ + \sup_{\xi'} \int_{|\xi'| \leq 2N, |\xi''| \geq 3N} |\mathbf{k}_w^{j,k}(V_j(s_1), \zeta', \xi'')| d\xi'' \quad (4.10) \\ \leq \frac{C_\eta}{\kappa^2} e^{-\frac{\eta}{8}N^2} \int_0^t \int_0^{s_1} e^{-\nu_0(t-s_1)} \|h(s_1)\|_{L^\infty} ds_1 ds \\ \leq C_\eta e^{-\frac{\eta}{8}N^2} e^{-\frac{\nu_0}{2}t} \sup_{0 \leq s \leq t} \{e^{\frac{\nu_0}{2}s} \|h(s)\|_{L^\infty}\}. \end{aligned}$$

CASE 3. $|\xi| \leq N$, $|\xi'| \leq 2N$, $|\xi''| \leq 3N$. This is the last remaining case because if $|\xi'| > 2N$, it is included in Case 2; while if $|\xi''| > 3N$, either $|\xi'| \leq 2N$ or $|\xi'| \geq 2N$ are also included in Case 2. We now can bound the second term in (4.6) by

$$C \int_0^t \int_B \int_0^s e^{-\nu_0(t-s_1)} |\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi') \mathbf{k}_w^{j,k}(V_j(s_1), \zeta', \xi'') h_k(s_1, X_j(s_1), \xi'')|$$

where $B = \{|\xi'| \leq 2N, |\xi''| \leq 3N\}$. We notice that $\mathbf{k}_w^{i,j}(\xi, \xi')$ has a possible integrable singularity of the type $\frac{1}{|\xi - \xi'|}$. We can choose $\mathbf{k}_N^{i,j}(\xi, \xi')$ smooth with compact support such that

$$\sup_{|p| \leq 3N} \int_{|\xi'| \leq 3N} |\mathbf{k}_N^{i,j}(p, \xi') - \mathbf{k}_w^{i,j}(p, \xi')| d\xi' \leq \frac{1}{N}. \quad (4.11)$$

Split $\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi') \mathbf{k}_w^{j,k}(V_j(s_1), \zeta', \xi'')$ into

$$\begin{aligned} & \{\mathbf{k}_w^{i,j}(V_i(s), \zeta, \xi') - \mathbf{k}_N^{i,j}(V_i(s), \zeta', \xi')\} \mathbf{k}_w^{j,k}(V_j(s_1), \zeta', \xi'') \\ & + \{\mathbf{k}_w^{i,j}(V_j(s_1), \zeta, \xi') - \mathbf{k}_N^{i,j}(V_j(s_1), \zeta, \xi')\} \mathbf{k}_N^{j,k}(V_j(s), \zeta', \xi'') \\ & + \mathbf{k}_N^{i,j}(V_i(s), \zeta, \xi') \mathbf{k}_N^{j,k}(V_j(s_1), \zeta', \xi''). \end{aligned}$$

We then integrate the first term above in ξ'' and the second term above in ξ' . By (4.8), we can use such an approximation (4.11) to bound the s_1, s integration by

$$\begin{aligned} & \frac{C e^{-\frac{\nu_0}{2}t}}{N} \sup_{0 \leq s \leq t} \{e^{\frac{\nu_0}{2}s} \|h(s)\|_{L^\infty}\} \\ & \times \left\{ \sup_{|\xi'| \leq 2N} \int |\mathbf{k}_w^{i,j}(V_j(s_1), \zeta', \xi'')| d\xi'' + \sup_{|\xi| \leq 2N} \int |\mathbf{k}_w^{j,k}(V_i(s), \zeta, \xi')| d\xi' \right\} \\ & + C \int_0^t \int_B \int_{s_1}^s e^{-\nu_0(t-s_1)} \mathbf{k}_N^{i,j}(V_i(s), \zeta, \xi') \mathbf{k}_N^{j,k}(V_j(s_1), \xi'') |h_j(s, X_j(s_1), \zeta', \xi'')|. \end{aligned} \quad (4.12)$$

The first term above is further bounded by $\frac{C e^{-\frac{\nu_0}{2}t}}{N} \sup_{0 \leq s \leq t} \{e^{\frac{\nu_0}{2}s} \|h(s)\|_{L^\infty}\}$.

Fix $\varepsilon > 0$. We use now Lemma 3.3 for the last main contribution in (4.13) for which we separate two cases $|X_j(s; t, x, v)| \geq L_\varepsilon$ and $|X_j(s; t, x, v)| \leq L_\varepsilon$ where L_ε is given in Lemma 3.3.

In the case $|X_j(s; t, x, v)| \geq L_\varepsilon$, we bound it by

$$\begin{aligned} & C_N \int_0^t \int_B \int_0^s e^{-\nu_0(t-s_1)} |h_k(s, X_j(s_1), \xi'')| \mathbf{1}_{|X_j(s_1)| \geq L_\varepsilon} ds d\xi'' d\xi' ds_1 \\ & \leq C_N \left\{ \int_0^{t-\varepsilon} + \int_{t-\varepsilon}^t \right\}. \end{aligned}$$

The second integral is bounded by $C \varepsilon e^{-\frac{\nu_0}{2}t} \sup_{0 \leq s \leq t} \{e^{\frac{\nu_0}{2}s} \|h(s)\|_{L^\infty}\}$. In the first integral, since $s \leq t - \varepsilon$, by Lemma 3.3, we can make a change of variable

$$y = X_j(s_1) = X_j(s_1; X_i(s; t, x, v), v') \quad (4.13)$$

because $|\frac{dy}{dv'}| \geq \frac{\varepsilon}{2}$. We observe since $\|\partial_x \phi\|_{L^\infty} \leq C$, that from (3.1)

$$\begin{aligned} |v' - V_j(\tau)| &\leq \int_\tau^s \|\partial_x \phi\|_{L^\infty} d\tau \leq T_0 \|\partial_x \phi\|_{L^\infty}, \\ |y - X_i(s)| &\leq \int_{s_1}^s |V_j(\tau)| d\tau \leq T_0(|v'| + T_0 \|\partial_x \phi\|_{L^\infty}) \leq C_{T_0, N} \end{aligned}$$

for $|v'| \leq 2N$. By first integrating over ζ' and using the change of variable (4.13)

$$\begin{aligned} &\int_0^{t-\varepsilon} \int_B \int_0^s e^{-\nu_0(t-s_1)} |h_k(s_1, X_j(s_1), \xi'')| \mathbf{1}_{|X_j(s_1)| \geq L_\varepsilon} ds d\xi'' d\xi' ds_1 \\ &\leq \frac{C}{\varepsilon} \int_0^{t-\varepsilon} \int_{|y-X_i(s)| \leq C_{T_0, N}} \int_{|\xi''| \leq 3N} \int_0^s e^{-\nu_0(t-s_1)} |h_k(s_1, y, \xi'')| ds d\xi'' dy ds_1 \\ &\leq \frac{C_N}{\varepsilon} \sup_{0 \leq s_1 \leq T_0} \int_{|y-X_i(s)| \leq C_{T_0, N}} \int_{|\xi''| \leq 3N} |h_k(s_1, y, \xi'')| ds_1 d\xi'' dy \\ &= \int_{|f_k(t)-M_k| \geq \kappa M_j} + \int_{|f_k(t)-M_k| \leq \kappa M_j} \\ &\leq C_{T_0, N, \varepsilon} \{ \mathcal{H}(g(0)) + \sqrt{\mathcal{H}(g(0))} \}. \end{aligned}$$

We have used the fact $h_k = \frac{w(f_k - M_k)}{\sqrt{M_k}}$, (which is bounded by $f_k - M_k$ for $|\xi''| \leq 3N$), and applied Lemma 2.2.

For $|X_i(s; t, x, v)| \leq L_\varepsilon$, for any $\eta > 0$, we again employ Lemma 3.3 to find O_{x_l} such that

$$\begin{aligned} &\int_0^{T_0} \int_B \int_0^{T_0} e^{-\nu_0(t-s_1)} |h_k(s_1, X_j(s_1), \xi'')| \mathbf{1}_{|X_j(s)| \leq L_\varepsilon} ds d\xi'' d\xi' ds_1 \\ &= \int_0^{T_0} \int_B \int_0^{T_0} \mathbf{1}_{O_{x_l}^c} e^{-\nu_0(t-s_1)} |h_k(s_1, X_j(s_1), \xi'')| \mathbf{1}_{|X_j(s)| \leq L_\varepsilon} ds d\xi'' d\xi' ds_1 \\ &+ \int_0^{T_0} \int_B \int_0^{T_0} \mathbf{1}_{O_{x_l}} e^{-\nu_0(t-s_1)} |h_k(s_1, X_j(s_1), \xi'')| \mathbf{1}_{|X_j(s)| \leq L_\varepsilon} ds d\xi'' d\xi' ds_1. \end{aligned}$$

Since $|[0, T_0] \times [-N, N] \cap O_{x_l}^c| < \eta$, the first part is bounded by

$$C_{T_0, N, \varepsilon} \eta e^{-\frac{\nu_0}{2} t} \sup_{0 \leq s \leq t} \{ e^{\frac{\nu_0}{2} s} \|h(s)\|_{L^\infty} \}.$$

The second part is bounded by

$$C_{T_0, N, \varepsilon} \int_0^{T_0} \int_0^{T_0} \int_B \mathbf{1}_{O_{x_l}} |h_k(s_1, X_j(s_1), \xi'')| ds ds_1 d\xi' d\xi''.$$

Since $|\frac{\partial X_j(s_1; X_i(s; t, x, v), v')}{\partial v'}| > m_\eta/2$ on O_{x_1} from Lemma 3.3, we can make a (local) change of variable $y = X_j(s_1) = X_j(s_1; X_i(s; t, x, v), v')$ to get

$$\begin{aligned} & C_{T_0, N, \varepsilon} \int_0^{T_0} \int_0^{T_0} \int_B \mathbf{1}_{O_{x_1}^c} |h_k(s_1, X_j(s_1), \xi'')| ds ds_1 d\xi' d\xi'' \\ &= C_{T_0, N, \varepsilon} \sum_{I_i} C_{T_0, N} \int_0^{T_0} \int_0^{T_0} \int_{|y - X_i(s_1)| \leq C_{T_0, N}} \times \\ & \quad \int_{|\xi''| \leq 3N} h_k(s_1, y, \xi'') ds ds_1 dy d\xi'' \\ &= \int_{|f_k(t) - M_k| \leq \kappa M_k} + \int_{|f_k(t) - M_k| \geq \kappa M_k} \\ &\leq C_{T_0, N, \varepsilon, \eta} \{\mathcal{H}(g(0)) + \sqrt{\mathcal{H}(g(0))}\}. \end{aligned}$$

Collecting terms, we conclude

$$\begin{aligned} \sup_{0 \leq s \leq T_0} e^{\frac{\nu_0}{2}t} \|h(s)\|_{L^\infty} &\leq C(1 + T_0) \|h(0)\|_{L^\infty} \\ &+ \left\{ \frac{C_{T_0}}{N} + C_{N, T_0} \varepsilon + C_{N, T_0, \varepsilon, \eta} \right\} \sup_{0 \leq s \leq T_0} \{e^{\frac{\nu_0}{2}s} \|h(s)\|_{L^\infty}\} \\ &+ C \left\{ \sup_{0 \leq s \leq T_0} e^{\frac{\nu_0}{2}s} \|h(s)\|_{L^\infty} \right\}^2 + C_{T_0, N, \varepsilon, \eta} \sqrt{\mathcal{H}(g(0))}. \end{aligned}$$

Assume $\sup_{0 \leq s \leq T_0} \|h(s)\|_{L^\infty}$ is sufficiently small. We first choose T_0 sufficiently large so that

$$2C(1 + T_0)e^{-\frac{\nu_0}{2}T_0} \leq \frac{1}{2},$$

then N sufficiently large, then ε sufficiently small, finally η small to conclude our lemma.

Proof of Theorem 1.2: Assume $\sup_{0 \leq t \leq \infty} \|h(t)\|_{L^\infty}$ is small. We first establish (1.24). Choose any $n = 0, 1, 2, 3, \dots$ and apply Lemma 4.1 repeatedly to get

$$\begin{aligned} \|h(nT_0)\|_{L^\infty} &\leq \frac{1}{2} \|h(\{n-1\}T_0)\|_{L^\infty} + C_{T_0} \sqrt{\mathcal{H}(g(0))} \\ &\leq \frac{1}{4} \|h(\{n-2\}T_0)\|_{L^\infty} + \frac{1}{2} C_{T_0} \sqrt{\mathcal{H}(g(0))} + C_{T_0} \sqrt{\mathcal{H}(g(0))} \\ &\leq \dots \\ &\leq \frac{1}{2^n} \|h_0\|_{L^\infty} + C_{T_0} \sqrt{\mathcal{H}(g(0))} \left\{ 1 + \frac{1}{2} + \frac{1}{4} + \dots \right\} \\ &\leq \frac{1}{2^n} \|h_0\|_{L^\infty} + 2C_{T_0} \sqrt{\mathcal{H}(g(0))}. \end{aligned}$$

For any t , we can find n such that $nT_0 \leq t \leq \{n+1\}T_0$, and from L^∞ estimate from $[0, T_0]$, we conclude (1.24) by

$$\|h(t)\|_{L^\infty} \leq C_{T_0} \|h(nT_0)\| \leq C \{ \|h_0\|_{L^\infty} + \sqrt{\mathcal{H}(g(0))} \}.$$

To prove (1.25), we take x and v derivatives to get

$$\begin{aligned} & \{\partial_t + v\partial_x + F(M_{i+1} + \sqrt{M_{i+1}g_{i+1}})\partial_v + \nu(\xi)\}\partial_x g_i - K^{i,j}\partial_x g_j \\ &= -\partial_x F(M_{i+1} + \sqrt{M_{i+1}g_{i+1}})\partial_v g_i + \beta\partial_x F(\sqrt{M_{i+1}g_{i+1}})v\partial_x \sqrt{M_i} + \\ & \quad + \partial_x \{F(\sqrt{M_j g_j})v g_i\} + \partial_x \{\Gamma(g_i, g_i) + \Gamma(g_i, g_j)\}; \end{aligned} \quad (4.14)$$

$$\begin{aligned} & \{\partial_t + v\partial_x + F(M_j + \sqrt{M_j g_j})\partial_v + \nu(\xi)\}\partial_v g_i - \partial_v \{K^{i,j} g_j\} + \{\partial_v \nu(\xi)\}g_i \\ &= -\partial_x g_i + \beta\partial_x F(\sqrt{M_j g_j})v\sqrt{M_i} + \beta\partial_x F(\sqrt{M_j g_j})v\partial_x \sqrt{M_i} \\ & \quad + F(\sqrt{M_j g_j})\partial_v \{v g_i\} + \partial_v \{\Gamma(g_i, g_i) + \Gamma(g_i, g_j)\} \end{aligned} \quad (4.15)$$

where $K^{i,j}$ has similar property as K_i in [6] (see Lemma 2.2 in [6], p. 1109. In particular, $\|\partial_v \{K^{i,j} g_j\} \partial_v g_i\|_{L^1} \leq \frac{1}{2} \|\partial_v g\|_\nu^2 + C \|g\|_{L^2}^2$ so that a positive dissipation for $\partial_v g_i$ occur for small $\|h\|_{L^\infty}$ in (4.15). Notice that $L = \nu - K \geq 0$. We take inner product with $\partial_x g_i$ and $\partial_v g_i$ respectively, following the procedures in [6] to get:

$$\begin{aligned} & \frac{d}{dt} \frac{1}{2} \|\partial_x g\|_{L^2}^2 \leq C \{\|\partial_x F(M_{i+1})\|_{L^\infty} + \|h\|_{L^\infty}\} \|\nabla_{x,v} g\|_{L^2}^2 + C \|g\|_{L^2}^2. \\ & \frac{d}{dt} \frac{1}{2} \|\partial_v g\|_{L^2}^2 + \frac{1}{4} \|\partial_v g\|_\nu^2 \leq C \|\partial_x g\|_{L^2}^2 + C \|g\|_{L^2}^2. \end{aligned} \quad (4.16)$$

Hence (1.25) follows from the Gronwall Lemma since $\sup_{0 \leq t \leq \infty} \|h(t)\|_{L^\infty}$ is bounded by (1.24). With such an estimate, we obtain the uniqueness by taking L^2 estimate for the difference for (1.17) because the most difficult term $F(M_{i+1} + \sqrt{M_{i+1}g_{i+1}})\partial_v g_i$ can be handled.

5 Linear Instability: Growing Mode.

In this section we study the linearization of the equation (1.1) around the homogeneous equilibrium $M_{\text{hom}} = (\mu_\beta, \mu_\beta)$. In the sequel we omit the index β for sake of shortness: $\mu = \mu_\beta = (\frac{\beta}{2\pi})^{3/2} e^{-\beta \frac{|\xi|^2}{2}}$. When M is replaced by $M_{\text{hom}} = (\mu, \mu)$ in (1.17), we get the following linearized Vlasov-Boltzmann system:

$$\partial_t g + \mathcal{L}g = 0, \quad (5.1)$$

where $g = (g_1, g_2)$,

$$(\mathcal{L}g)_i = v\partial_x g_i - \beta F(\sqrt{\mu}g_{i+1})v\sqrt{\mu} - L_i g$$

and

$$L_i g = \frac{1}{\sqrt{\mu}} \left(Q(\sqrt{\mu}g_i, 2\mu) + Q(\mu, \sqrt{\mu}(g_1 + g_2)) \right).$$

We seek an exponential growing mode for such a system when $\beta > 1$. To this end, we consider a family of systems

$$\partial_t g + \mathcal{L}^\alpha g = 0, \quad (5.2)$$

$$(\mathcal{L}^\alpha g)_i = v\partial_x g_i - \beta F(\sqrt{\mu}g_{i+1})v\sqrt{\mu} - \alpha L_i g$$

and show that there is a growing mode for all $\alpha > 0$. We seek a growing mode periodic in x , so we assume periodic dependence on space and exponential in time: $g_1(t, x, v, \zeta) = e^{\lambda t} e^{ikx} q(v, \zeta)$, $g_2(t, x, v, \zeta) = e^{\lambda t} e^{-ikx} q(-v, \zeta)$, so that the system (5.2), using the definition of F (see (1.2)), reduces to the single equation

$$\{\lambda + \mathbf{i}vk\}q - \beta k \mathbf{i} \hat{U}(k) \left\{ \int q \sqrt{\mu} d\xi \right\} v \sqrt{\mu} = \alpha Lq, \quad (5.3)$$

with $Lg = \frac{2}{\sqrt{\mu}} \left(Q(\sqrt{\mu}g, \mu) + Q(\mu, \sqrt{\mu}g) \right)$. Equivalently, $q(\xi)$ is an eigenfunction for the operator \mathcal{T}^α with eigenvalue λ :

$$(\mathcal{T}^\alpha q)(\xi) = \mathbf{i}vkq(\xi) - \beta k \mathbf{i} \hat{U}(k) \left\{ \int q(\xi) \sqrt{\mu} d\xi \right\} v \sqrt{\mu} - \alpha Lq(\xi). \quad (5.4)$$

Lemma 5.1 *Let $\beta > 1$. There exists sufficiently small $\alpha > 0$ such that there is an eigenfunction $q(\xi)$ to \mathcal{T}^α with $\text{Re}\lambda > 0$.*

Proof We first study the eigenvalue problem for the unperturbed operator \mathcal{T}^0 for $\alpha = 0$:

$$(\lambda + \mathbf{i}vk)q - \beta k \mathbf{i} \hat{U}(k) \left\{ \int_{\mathbb{R}^3} q \sqrt{\mu} d\xi \right\} v \sqrt{\mu} = 0. \quad (5.5)$$

This is similar to the Penrose dispersion relation for the unperturbed Vlasov-Poisson system ([14]) for a collisionless plasma. From (5.5), we obtain:

$$q = \frac{\beta k \hat{U}(k) \mathbf{i} \left\{ \int_{\mathbb{R}^3} q \sqrt{\mu} d\xi \right\} v \sqrt{\mu}}{(\lambda + \mathbf{i}vk)}. \quad (5.6)$$

Normalizing $\int_{\mathbb{R}^3} q \sqrt{\mu} d\xi = 1$, we deduce that

$$\beta \mathbf{i} \int_{\mathbb{R}^3} \frac{\hat{U}(k) k v \mu(\xi)}{(\lambda + \mathbf{i}vk)} d\xi = 1.$$

Multiply and divide (5.6) by $(\lambda - \mathbf{i}vk)\sqrt{\mu}$, take the imaginary part and then integrate on ξ . By consistency we must have:

$$\beta \int_{\mathbb{R}^3} \frac{v^2 \hat{U}(k) k^2 \mu(\xi)}{\lambda^2 + k^2 v^2} d\xi = 1.$$

Define $F(\lambda, k) \equiv \beta \int_{\mathbb{R}^3} \frac{v^2 \hat{U}(k) k^2 \mu(\xi)}{\lambda^2 + k^2 v^2} d\xi$. Clearly $F(0, k) = \beta \hat{U}(k)$. Since $\hat{U}(0) = 1$ and $\beta > 1$, there is k_0 sufficiently small so that

$$F(0, k_0) = \beta \hat{U}(k_0) > 1. \quad (5.7)$$

Moreover, $\lim_{\lambda \rightarrow \infty} F(\lambda, k_0) = 0$ for any $k_0 \neq 0$. Hence there exists a real number (eigenvalue) $\lambda > 0$ such that $F(\lambda, k_0) = 1$ and $q = \frac{\beta k \hat{U}(k) \mathbf{i} v \sqrt{\mu}}{(\lambda + \mathbf{i}vk)}$ is the eigenfunction.

We now fix $k = k_0$ and return to (5.3).

It can be proved (see for example [9]) that $\|Lq\|_{L^2} \leq C\|\nu q\|_{L^2}$. Moreover, for hard spheres, there are constants C_1 and $C_2 > 0$ such that

$$\begin{aligned} \|\mathcal{T}^0 q\|_{L^2} &= \|\mathbf{i}k_0 \nu q - \beta k_0 \mathbf{i}\hat{U}(k_0) \left\{ \int_{\mathbb{R}^3} q \sqrt{\mu} d\xi \right\} \nu \sqrt{\mu}\|_{L^2} \\ &\geq C_1 \|\nu q\|_{L^2} - C_2 \|q\|_{L^2}. \end{aligned} \quad (5.8)$$

This implies $\|Lq\|_{L^2} \leq C\{\|\mathcal{T}^0 q\|_{L^2} + \|q\|_{L^2}\}$ so the perturbation L is \mathcal{T}^0 -bounded. Since

$$\mathcal{T}^\alpha = \mathcal{T}^0 - \alpha L,$$

we deduce from Kato's book (P206, [12]) that, for α small, there is an eigenvalue λ and an eigenfunction $q(v, \zeta)$ for (5.3) with positive real part λ for (5.3).

Theorem 5.1 *Let $\beta > 1$. Then there is a $\frac{2\pi}{k_0}$ -periodic eigenvector (\tilde{g}, \bar{g}) for $-\mathcal{L}$ such that $\operatorname{Re}\lambda > 0$ and $\tilde{g}(x, v, \zeta) = \tilde{g}(-x, -v, \zeta)$.*

Proof We fix k_0 as in Lemma 5.1. The idea is to define, for the family of equations (5.2)

$$\alpha_0 = \sup_{\alpha} \left\{ \alpha : \text{there is an eigenvalue with positive real part for } \mathcal{T}^\alpha \right. \\ \left. \text{with a } \frac{2\pi}{k_0} \text{-periodic eigenvector} \right\}.$$

By Lemma 5.1, for α sufficiently small this set is not empty. We want to show that $\alpha_0 = +\infty$.

We prove it by contradiction. Suppose $\alpha_0 < +\infty$. We claim that, if there is such a finite $\alpha_0 > 0$, then there is an eigenvalue λ_0 with an eigenfunction q_0 with $\|q_0\|_\nu = \sqrt{\int \nu |q_0|^2 d\xi} = 1$, such that $\operatorname{Re}\lambda_0 = 0$ and

$$(\lambda_0 + \mathbf{i}vk_0)q_0 - \beta k_0 \mathbf{i}\hat{U}(k_0) \left\{ \int_{\mathbb{R}^3} q_0 \sqrt{\mu} d\xi \right\} \nu \sqrt{\mu} = \alpha_0 Lq_0. \quad (5.9)$$

Proof of the claim: In fact, by (5.8), choose a family of eigenfunctions $q_\alpha \in L^2$ such that $\|q_\alpha\|_\nu = 1$, $\operatorname{Re}\lambda_\alpha > 0$, as $\alpha \rightarrow \alpha_0$ and

$$(\lambda_\alpha + \mathbf{i}vk_0)q_\alpha - \beta k_0 \mathbf{i}\hat{U}(k_0) \left\{ \int_{\mathbb{R}^3} q_\alpha \sqrt{\mu} d\xi \right\} \nu \sqrt{\mu} = \alpha Lq_\alpha. \quad (5.10)$$

Notice that both $(Lq_\alpha, \bar{q}_\alpha)$ and $(\mathbf{i}vk_0 q_\alpha, \bar{q}_\alpha)$ are bounded by $C\|q_\alpha\|_\nu^2$. As $\alpha \rightarrow \alpha_0$, taking L^2 inner product with \bar{q}_α for (5.10), we deduce that $|\lambda_\alpha|$ is bounded for $\alpha \rightarrow \alpha_0$. Hence $\lim_{\alpha \rightarrow \alpha_0} \lambda_\alpha = \lambda_0$ (up to subsequences) with $\operatorname{Re}\lambda_0 \geq 0$. We now prove that λ_0 is an eigenvalue so that $\operatorname{Re}\lambda_0 = 0$ by the definition of α_0 and the claim is proven. Clearly, we may assume that $\lim q_\alpha = q_0$ weakly in L^2 and (5.9) are valid as $\alpha \rightarrow \alpha_0$. We only need to show that $\lim q_\alpha = q_0$ strongly so that $\|q_0\|_\nu = 1$, and q_0 is an eigenfunction. Denote $\mathbf{P}h = \{1, v, |v|^2\}\sqrt{\mu}$. Clearly $\mathbf{P}q_\alpha \rightarrow \mathbf{P}q_0$ strongly in L^2 . It thus is

left to show that $(\mathbf{I} - \mathbf{P})q_\alpha \rightarrow (\mathbf{I} - \mathbf{P})q_0$ strongly in L^2 . We subtract (5.10) from (5.9) to get

$$\begin{aligned} & (\lambda_\alpha - \lambda_0)q_\alpha + (\alpha - \alpha_0)Lq_\alpha + (\lambda_0 + \mathbf{i}k_0v)(q_\alpha - q_0) \\ & - \beta k_0 \mathbf{i} \hat{U}(k_0) \left\{ \int_{\mathbb{R}^3} \{q_\alpha - q_0\} \sqrt{\mu} d\xi \right\} v \sqrt{\mu} + \alpha_0 L(q_\alpha - q_0) \\ & = 0. \end{aligned}$$

We take the L^2 inner product with $\bar{q}_\alpha - \bar{q}_0$ and then take the real part. Since $\operatorname{Re} \lambda_0 \|q_\alpha - q_0\|^2 \geq 0$, and $\mathbf{i}k_0 \int v |q_\alpha - q_0|^2 d\xi$ is purely imaginary, we obtain:

$$\begin{aligned} \alpha_0 \langle L(g^\alpha - g), (\bar{g}^\alpha - \bar{g}) \rangle & \leq (|\lambda_\alpha - \lambda_0| + |\alpha - \alpha_0|) \|q_\alpha\|_\nu \cdot \|q_\alpha - q_0\|_\nu \\ & + C \left| \int_{\mathbb{R}^3} \{q_\alpha - q_0\} \sqrt{\mu} d\xi \right| \cdot \|(g^\alpha - g)\|. \end{aligned}$$

Therefore, $(\mathbf{I} - \mathbf{P})\{g^\alpha - g\} \rightarrow 0$ in L^2_ν and $\|g\|_\nu = 1$ and our claim follows.

Hence, λ_0 is purely imaginary. Actually we show that it is 0. To do this we take inner product with \bar{q}_0 in (5.9) to get

$$\lambda_0 \|q_0\|_2^2 - \beta k_0 \mathbf{i} \hat{U}(k_0) \left\{ \int_{\mathbb{R}^3} q_0 \sqrt{\mu} d\xi \right\} \left\{ \int_{\mathbb{R}^3} \bar{q}_0 v \sqrt{\mu} d\xi \right\} + \alpha_0 \langle Lg, \bar{g} \rangle = 0.$$

But by integrating $\sqrt{\mu} \times (5.9)$ over ξ , we obtain the continuity equation

$$\lambda_0 \int_{\mathbb{R}^3} q_0 v \sqrt{\mu} d\xi + k_0 \mathbf{i} \int_{\mathbb{R}^3} q_0 v \sqrt{\mu} d\xi = 0.$$

Therefore, $k_0 \mathbf{i} \int_{\mathbb{R}^3} \bar{q}_0 v \sqrt{\mu} d\xi = \bar{\lambda}_0 \int_{\mathbb{R}^3} \bar{q}_0 v \sqrt{\mu} d\xi$ and

$$\lambda_0 \|g\|_2^2 - \beta \bar{\lambda}_0 \hat{U}(k_0) \left| \int_{\mathbb{R}^3} q_0 \sqrt{\mu} d\xi \right|^2 + \alpha_0 \langle Lg, \bar{g} \rangle = 0. \quad (5.11)$$

Since $\hat{U}(k_0)$ is real and λ_0 is purely imaginary, taking the real part of (5.11) we conclude that $\alpha_0 \langle Lq_0, \bar{q}_0 \rangle = 0$. Therefore, q_0 is a linear combination of the collision invariants $\sqrt{\mu}$, $\xi \sqrt{\mu}$ and $|\xi^2|^2 \sqrt{\mu}$ and $\alpha_0 Lq_0$ vanishes in (5.9). Now (5.9) reduces to a pure Vlasov equation, and we deduce that

$$q_0(\xi) = \frac{\beta \mathbf{i} k_0 \hat{U}(k_0) \left\{ \int q_0(\xi) \sqrt{\mu} d\xi \right\} v \sqrt{\mu}}{\lambda_0 + \mathbf{i} v j k_0}.$$

Since $\int q_0(\xi) \sqrt{\mu} d\xi \neq 0$, this is compatible with the condition that q_0 is a combination of collision invariants if and only if $\lambda_0 = 0$. Thus $q_0(\xi) = \beta \hat{U}(k_0) \sqrt{\mu}$ and hence

$$\beta \hat{U}(k_0) = 1 \quad (5.12)$$

which is a contradiction to (5.7). Theorem 5.1 thus follows.

6 Nonlinear Instability.

In order to establish the non linear instability, we need several lemmas on the properties of the fastest linear growing mode. First of all we need to establish the smoothness and long time behavior for the growing mode. Recall the definition of the operator \mathcal{L} (5.1). We define

$$\mathcal{M} = \{g = [g_1, g_2] \in L^2 \mid g_1(x, v, \zeta) = g_2(-x, -v, \zeta)\}$$

and $\|\cdot\|_{L^2(\mathcal{M})}$ will denote the L^2 norm on this set. We have the following lemmas:

Lemma 6.1 *Let $\beta > 1$. Then for k_0 sufficiently small, for all $\delta > 0$, the spectrum of $-\mathcal{L}$ in $\{\operatorname{Re}\lambda > \delta\}$ consists of a finite number of eigenvalues of finite multiplicity. If λ_1 denotes an eigenvalue with maximal real part, and $\Lambda > \max\{0, \operatorname{Re}\lambda_1\}$, then there exists $C_\Lambda > 0$ such that, for any $g_0 \in \mathcal{M}$,*

$$\|e^{-t\mathcal{L}}g_0\|_{L^2(\mathcal{M})} \leq C_\Lambda e^{\Lambda t} \|g_0\|_{L^2(\mathcal{M})}.$$

Proof This follows easily from the Vidav's Lemma [16]. Notice that we can split

$$\begin{aligned} \mathcal{L}g &= \{v\partial_x g + Lg\} + \{\beta F(\sqrt{\mu}g)v\sqrt{\mu}\} \\ &\equiv \mathcal{A}g + \mathcal{K}g. \end{aligned}$$

where \mathcal{K} is a compact operator from L^2 to L^2 , while $\|e^{-t\mathcal{A}}\|_{L^2(\mathcal{M}) \rightarrow L^2(\mathcal{M})} \leq 1$.

Lemma 6.2 *Let $R = (R_1, R_2) \in L^2(\mathcal{M})$ with $\|R\|_{L^2(\mathcal{M})} = 1$ be an eigenvector of $-\mathcal{L}$ with $\operatorname{Re}\lambda > 0$. Then there exists a constant C depending only on λ such that*

$$\|\nabla_{x,v}R\|_{L^2(\mathcal{M})} \leq C. \quad (6.1)$$

$$\sup_{x,v} w(v)|R(x, v)| \leq C, \quad (6.2)$$

where w is a polynomial weight as in previous section.

Proof We begin with $R \in L^2$. We first claim

$$R = - \int_0^\infty e^{-\lambda t} e^{-t\mathcal{A}} \mathcal{K}R dt. \quad (6.3)$$

Notice that the corresponding growing mode $g(t) = e^{\lambda t}R$ satisfies

$$\partial_t g + \mathcal{A}g = -\mathcal{K}g$$

so that $e^{\lambda t}R = e^{-(t-s)\mathcal{A}}e^{\lambda s}R - \int_s^t e^{-(t-\tau)\mathcal{A}}\mathcal{K}R d\tau$. Letting $s \rightarrow -\infty$, since $\|e^{-t\mathcal{A}}\|_{L^2 \rightarrow L^2} \leq 1$ for any $t > 0$ and $\operatorname{Re}\lambda > 0$, we get

$$e^{\lambda t}R = - \int_{-\infty}^t e^{-(t-\tau)\mathcal{A}}\mathcal{K}R e^{\lambda\tau} d\tau = - \int_0^\infty e^{-\tau\mathcal{A}}\mathcal{K}R e^{\lambda(t-\tau)} d\tau.$$

Dividing by $e^{\lambda t}$ we prove our claim because $\operatorname{Re}\lambda > 0$ and the integral converges in L^2 .

From the property of linear Boltzmann equation, clearly $\|\partial_x\{e^{-t\mathcal{A}}g\}\|_{L^2(\mathcal{M})} \leq \|\partial_x g\|_{L^2(\mathcal{M})}$. Taking v derivative of $(\partial_t + v\partial_x + L)g = 0$ yields:

$$\{\partial_t + v\partial_x\}\{\partial_v g\} + \partial_v\{-Lg\} = -\partial_x g.$$

From [6], $\int \partial_v\{-Lg\}\partial_v g \geq \int \frac{\nu}{2}|\partial_v g|^2 - C_\nu\|g\|_{L^2}^2$. We thus obtain by taking L^2 inner product with $\partial_v g$,

$$\|\partial_v g\|_{L^2} = \|\partial_v\{e^{-t\mathcal{A}}g_0\}\|_{L^2} \leq C(t+1)\{\|g_0\|_{L^2} + \|\nabla_{x,v}g_0\|_{L^2}\}. \quad (6.4)$$

Since $\mathcal{K}R \in C^\infty$ and $\|\partial_x\{\mathcal{K}R\}\|_{L^2} + \|\partial_v\{\mathcal{K}R\}\|_{L^2} \leq C\|R\|_{L^2}$, we can take ∂_x and ∂_v derivatives in (6.3) to get

$$\begin{aligned} \|\partial_x R\|_{L^2} + \|\partial_v R\|_{L^2} &\leq \int_0^\infty e^{-\operatorname{Re}\lambda t} \{\|\partial_x\{e^{-t\mathcal{A}}\mathcal{K}R\}\|_{L^2} + \|\partial_v\{e^{-t\mathcal{A}}\mathcal{K}R\}\|_{L^2}\} dt \\ &\leq C \int_0^\infty e^{-\operatorname{Re}\lambda t} (t+1) \|\partial_x\{\mathcal{K}R\}\|_{L^2} + \|\partial_v\{\mathcal{K}R\}\|_{L^2} dt \\ &\leq C\|R\|_{L^2} \int_0^\infty e^{-\operatorname{Re}\lambda t} (t+1) dt \leq C\|R\|_{L^2}. \end{aligned}$$

We therefore deduce (6.1). To show (6.2), we denote $S = wR$. We then have

$$\lambda S = \{v\partial_x S + wL(\frac{S}{w})\} + \{\beta F(\frac{\sqrt{\mu}S}{w})vw\sqrt{\mu}\} \equiv \mathcal{A}_w S + \mathcal{K}_w S.$$

Applying the same proof in Section 3 for the stability for the pure linear Boltzmann operator, we can establish:

$$\|e^{-t\mathcal{A}_w}g_0\|_{L^\infty} \leq C\{\|g_0\|_{L^2} + \|g_0\|_{L^\infty}\}.$$

We can similarly obtain $S = -\int_0^\infty e^{-\lambda t} e^{-t\mathcal{A}_w} \mathcal{K}_w S dt$, so that from $\operatorname{Re}\lambda > 0$,

$$\begin{aligned} \|S\|_{L^\infty} &\leq \int_0^\infty e^{-\operatorname{Re}\lambda t} \|e^{-t\mathcal{A}_w} \mathcal{K}_w S\|_{L^\infty} dt \\ &\leq C \int_0^\infty e^{-\operatorname{Re}\lambda t} \{\|\mathcal{K}_w S\|_{L^\infty} + \|\mathcal{K}_w S\|_{L^2}\} dt \\ &\leq C. \end{aligned}$$

Lemma 6.3 *Let R be an eigenvector of $-\mathcal{L}$ with its eigenvalue λ with $\operatorname{Re}\lambda > 0$. If λ is not real, then there is a constant $\zeta > 0$ such that for all $t > 0$,*

$$\|e^{-\mathcal{L}t} \operatorname{Im}R\|_{L^2} \geq \zeta e^{\operatorname{Re}\lambda t} \|\operatorname{Im}R\|_2 > 0.$$

Proof We prove by contradiction. Notice that, since $\operatorname{Im}F(\sqrt{\mu}R) = F(\sqrt{\mu}\operatorname{Im}R)$, one can immediately check that

$$e^{-\mathcal{L}t} \operatorname{Im}R = \operatorname{Im}\{e^{-\mathcal{L}t}R\} = e^{\operatorname{Re}\lambda t} (\sin[\operatorname{Im}\lambda t] \operatorname{Re}R + \cos[\operatorname{Im}\lambda t] \operatorname{Im}R).$$

If the Lemma were false, by passing through a convergent subsequence of $\sin[\operatorname{Im}\lambda t_n]$ and $\cos[\operatorname{Im}\lambda t_n]$, with $n \rightarrow \infty$ we would have $a\operatorname{Im}R + b\operatorname{Re}R = 0$, with $a^2 + b^2 = 1$. Therefore either $\operatorname{Im}R$ or $\operatorname{Re}R$ would be a real eigenvector and λ would be real, a contradiction.

Lemma 6.4 *Let R be as in the preceding lemma. There exists $\delta_0 > 0$, such that for $0 < \delta < \delta_0$, there exists a (compactly supported) approximate eigenfunction R_δ such that*

$$\begin{aligned} \delta |R_\delta(x, v)|\sqrt{\mu} &\leq \mu, \\ \|R - R_\delta\|_{L^2} &\leq \sqrt{\delta}, \\ \|\partial_x R_\delta\|_{L^2} + \|\partial_v R_\delta\|_{L^2} &\leq C\{\|\partial_x R\|_{L^2} + \|\partial_v R\|_{L^2}\}. \end{aligned}$$

Proof In fact, we choose $\chi(v)$ be a smooth cutoff function $\chi(v) = 1$ for $|v| \leq N$ and $\chi(v) \equiv 0$ for $|v| \geq N + 1$. By Lemma 6.2, we have

$$|\chi(v)R(x, v)|\sqrt{\mu} \leq |R(x, v)|\sqrt{\mu} = |wS|\sqrt{\mu} \leq \{Cw\mu^{-1/2}\}\mu. \quad (6.5)$$

Define N by the equation $\delta = \frac{\mu^{1/2}(N+1)}{Cw(N+1)}$ and define

$$R_\delta = \chi(v)R(x, v).$$

Clearly the third estimate in the Lemma is valid. From (6.5) and the definition of N and δ , the first inequality in the Lemma is also valid. Since w is a polynomial, we have $\sqrt{\delta} = \sqrt{\frac{\mu^{1/2}(N+1)}{Cw(N+1)}} \geq \sqrt{\mu(N)}$ when N is large. We then conclude the lemma by

$$\|R - R_\delta\|_{L^2} = \|R\mathbf{1}_{|v| \geq N}\|_{L^2} \leq C \int_{|v| \geq N} \frac{\mu^{\frac{1}{2}}(v)}{w(v)} dv = C\mu^{\frac{1}{2}}(N) \leq \sqrt{\delta}.$$

We now establish the crucial bootstrap Lemma which shows that L^2 growth leads to the same growth rate for L^∞ .

Lemma 6.5 *Let $g = (g_1, g_2)$ be a solution to the nonlinear problem around M_{hom} :*

$$\begin{aligned} (\partial_t + v\partial_x)g_i + \beta F(\sqrt{\mu}g_{i+1})\sqrt{\mu}v - L_i g & \quad (6.6) \\ = -F(\sqrt{\mu}g_{i+1})\partial_v g_i + F(\sqrt{\mu}g_{i+1})g_i v + \Gamma(g_i, g_i) + \Gamma(g_i, g_{i+1}). & \end{aligned}$$

Assume that $\text{Re}\lambda > 0$ and

$$\|g(t)\|_{L^2} \leq Ce^{\text{Re}\lambda t} \|g(0)\|_{L^2}$$

for $t \in [0, T]$. There exists $\varepsilon_0 > 0$ such that if $\sup_{0 \leq t \leq T} \{\|wg(t)\|_{L^\infty} + \|g(t)\|_{L^2}\} \leq \varepsilon_0$, then there is a constant C such that

$$\begin{aligned} \|\partial_x g(t)\|_{L^2} + \|\partial_v g(t)\|_{L^2} + \|wg(t)\|_{L^\infty} \\ \leq Ce^{\text{Re}\lambda t} \{\|\partial_x g(0)\|_{L^2} + \|\partial_v g(0)\|_{L^2} + \|h(0)\|_{L^\infty}\}. \end{aligned} \quad (6.7)$$

Proof We take x and v derivatives for (6.6). Since $F(M_i) = 0$, from (4.16), we have ($h = wg$)

$$\begin{aligned} \frac{d}{dt} \|\partial_x g\|_{L^2}^2 &\leq C \|h\|_{L^\infty} \{\|\partial_x g\|_{L^2}^2 + \|\partial_v g\|_{L^2}^2\} + C \|g\|_{L^2}^2, \\ \frac{d}{dt} \|\partial_v g\|_{L^2}^2 &\leq C \|\partial_x g\|_{L^2}^2 + C \|g\|_{L^2}^2. \end{aligned} \quad (6.8)$$

Applying the Gronwall's inequality to (6.8), by $\|h\|_{L^\infty} \leq \varepsilon_0 < \operatorname{Re}\lambda$, we obtain

$$\begin{aligned} \|\partial_x g(t)\|_{L^2}^2 &\leq \|\partial_x g(0)\|_{L^2}^2 + C\varepsilon_0 \int_0^t e^{C\varepsilon_0(t-s)} \|\partial_v g(s)\|_{L^2}^2 ds \\ &\quad + C e^{\operatorname{Re}\lambda t} \|g(0)\|_{L^2}^2, \end{aligned} \quad (6.9)$$

$$\begin{aligned} \frac{d}{dt} \|\partial_v g(t)\|_{L^2}^2 &\leq C\varepsilon_0 \int_0^t e^{C\varepsilon_0(t-s)} \|\partial_v g(s)\|_{L^2}^2 ds + C e^{\operatorname{Re}\lambda t} \{\|g(0)\|_{L^2}^2 \\ &\quad + \|\partial_x g(0)\|_{L^2}^2\}. \end{aligned} \quad (6.10)$$

Letting $Y(t) = \int_0^t e^{C\varepsilon_0(t-s)} \|\partial_v g(s)\|_{L^2}^2 ds$ so that $Y' = \|\partial_v g(s)\|_{L^2}^2$. Multiplying with $Y' \geq 0$, we have

$$\frac{d}{dt} \{Y'\}^2 \leq C\varepsilon_0 \{Y'\}^2 + C\varepsilon_0 e^{\operatorname{Re}\lambda t} \{\|g(0)\|_{L^2}^2 + \|\partial_x g(0)\|_{L^2}^2\}.$$

Gronwall's inequality and integration by parts implies

$$\begin{aligned} \{Y'\}^2(t) &\leq \|\partial_v g(0)\|_{L^2}^4 + C \{\|g(0)\|_{L^2}^2 + \|\partial_x g(0)\|_{L^2}^2\}^2 \int_0^t e^{C\varepsilon_0(t-s)} e^{2\operatorname{Re}\lambda s} ds \\ &\leq C \{\|\partial_v g(0)\|_{L^2}^2 + \|g(0)\|_{L^2}^2 + \|\partial_x g(0)\|_{L^2}^2\}^2 e^{2\operatorname{Re}\lambda t}. \end{aligned}$$

Taking the square root and plugging into (6.9), we thus conclude the desired estimate for $\|\partial_v g(s)\|_{L^2}^2 + \|\partial_x g(t)\|_{L^2}^2$ in (6.7).

We now consider the L^∞ estimate. By repeating the same (but simpler) argument in Lemma 4.1 for the full nonlinear system (with $F(M_i) = 0$), we can show that if $\{\|h(t)\|_{L^\infty} + \|g(t)\|_{L^2}\}$ is sufficiently small,

$$\|h(t)\|_{L^\infty} \leq C \{\|h(0)\|_{L^\infty} + \sup_{0 \leq s \leq t} \|g(s)\|_{L^2}\} \leq C e^{\operatorname{Re}\lambda t} \{\|h(0)\|_{L^\infty} + \|g(0)\|_{L^2}\}.$$

Proof (Proof of Theorem 1.4) We choose R to be the eigenfunction whose eigenvalue has the largest positive real part. If λ is not real, then $\|\operatorname{Im}R\|_{L^2} = r > 0$. We choose the approximate eigenfunction $\operatorname{Im}R^\delta$ to the imaginary part of R by Lemma 6.4. In case λ is real, we simply do not take the imaginary parts.

We choose a family of solutions $f^\delta(0, x, \xi) = \mu + \delta \operatorname{Im}R^\delta \sqrt{\mu} \geq 0$ or $g^\delta(0, x, \xi) = \delta \operatorname{Im}R^\delta$. Note that the positivity follows from the first statement in Lemma 6.4, for δ sufficiently small. Clearly, from Lemma 6.4,

$$\|g^\delta(0) - \delta \operatorname{Im}R\|_{L^2} = \delta \|R - \operatorname{Im}R^\delta\|_{L^2} \leq \delta^{1+\frac{1}{2}} \leq \frac{\delta r}{2},$$

for δ sufficiently small. Hence, from Lemma 6.4,

$$\|g^\delta(0)\|_{H^1} + \|h^\delta(0)\|_{L^\infty} = \delta\|\text{Im}R^\delta\|_{H^1} + \delta\|\text{Im}R^\delta\|_{L^\infty} \leq C\delta r.$$

Now from the nonlinear Vlasov-Boltzmann system (6.6), we have

$$\begin{aligned} g^\delta(t) &= \delta e^{-\mathcal{L}t} \text{Im}R^\delta \\ &+ \int_0^t e^{-\mathcal{L}(t-\tau)} \left(-F(\sqrt{\mu}g_2)\partial_v g_1 + F(\sqrt{\mu}g_2)v g_1 + \Gamma(g_1, g_1) + \Gamma(g_1, g_2) \right. \\ &\quad \left. - F(\sqrt{\mu}g_1)\partial_v g_2 + F(\sqrt{\mu}g_1)v g_2 + \Gamma(g_2, g_2) + \Gamma(g_2, g_1) \right) d\tau. \end{aligned} \quad (6.11)$$

We choose Λ such that

$$\text{Re}\lambda < \Lambda < \left(1 + \frac{1}{2}\right)\text{Re}\lambda. \quad (6.12)$$

Let

$$T^{\delta\delta} = \frac{1}{\Lambda - \text{Re}\lambda} \left| \ln \frac{\zeta r}{2C_A\sqrt{\delta}} \right|,$$

By (6.12) and $T^\delta = \frac{1}{\text{Re}\lambda} \ln \frac{\theta}{\delta}$. Since $2(\Lambda - \text{Re}\lambda) < \text{Re}\lambda$, for small δ , we have $T^\delta \leq T^{\delta\delta}$. Clearly, $e^{(\Lambda - \text{Re}\lambda)t} \leq e^{(\Lambda - \text{Re}\lambda)T^{\delta\delta}} = \frac{\zeta r}{2C_A\sqrt{\delta}}$, and for $0 \leq t \leq T^{\delta\delta}$:

$$C_A\sqrt{\delta}e^{At} \leq \frac{\zeta}{2}e^{\text{Re}\lambda t}r. \quad (6.13)$$

Let

$$T^* = \sup_s \{s : \|\nabla_x g^\delta(t)\|_{L^2} + \|\nabla_v g^\delta(t)\|_{L^2} + \|h^\delta(t)\|_{L^\infty} \leq \varepsilon_0\} \quad (6.14)$$

$$T^{**} = \sup_s \{s : \|g^\delta(t) - \delta e^{-\mathcal{L}t}R^\delta\|_{L^2} \leq \frac{\zeta}{4}\delta e^{\text{Re}\lambda t}r, \text{ for all } 0 \leq t \leq s\} \quad (6.15)$$

For $0 \leq t \leq \min\{T^\delta, T^{**}\}$, we have from (6.13),

$$\begin{aligned} \|g^\delta(t)\|_{L^2} &\leq \delta\|e^{-\mathcal{L}t}\text{Im}R^\delta\|_{L^2} + \frac{\zeta}{4}\delta e^{\text{Re}\lambda t}r \\ &= \delta\|e^{-\mathcal{L}t}\text{Im}R\|_{L^2} + \delta\|e^{-\mathcal{L}t}\{R - R^\delta\}\|_{L^2} + \frac{\zeta}{4}\delta e^{\text{Re}\lambda t}r \\ &\leq \delta e^{\text{Re}\lambda t} + C_A\delta^{\{1+\frac{1}{2}\}}e^{At} + \frac{\zeta}{4}\delta e^{\text{Re}\lambda t}r \\ &\leq \left(1 + \frac{3\zeta}{4}\right)e^{\text{Re}\lambda t}\|g^\delta(0)\|_{L^2}. \end{aligned} \quad (6.16)$$

We now claim that $T^\delta \leq \min\{T^*, T^{**}\}$.

In fact, if $T^* < \min\{T^\delta, T^{**}\}$, then by (1.26), (6.14), (6.16), and the Lemma 6.5 (bootstrap lemma), we obtain for $0 \leq t \leq T^*$:

$$\|\nabla_{x,v}g^\delta(t)\|_{L^2} + \|h^\delta(t)\|_{L^\infty} \leq C e^{\text{Re}\lambda t} \{\|\nabla_{x,v}g^\delta(0)\|_{L^2} + \|h^\delta(0)\|_{L^\infty}\}.$$

In particular, by (1.26),

$$\begin{aligned} \|\nabla_{x,v}g^\delta(T^*)\|_{L^2} + \|h^\delta(T^*)\|_{L^\infty} &\leq C e^{\operatorname{Re}\lambda T^*} \{\|\nabla_{x,v}g^\delta(0)\|_{L^2} + \|h^\delta(0)\|_{L^\infty}\} \\ &< C e^{\operatorname{Re}\lambda T^\delta} \delta = C\theta. \end{aligned}$$

This is a contradiction to the definition of T^* if θ is chosen $\ll \varepsilon_0$.

On the other hand, if $T^{**} < \min\{T^\delta, T^*\}$, then by (6.12), (6.16), and Lemma 6.5,

$$\begin{aligned} \|g^\delta(T^{**}) - \delta e^{-\mathcal{L}T^{**}} \operatorname{Im}R^\delta\|_{L^2} &\leq C \int_0^{T^{**}} \{\|\nabla_{x,v}g^\delta(t)\|_{L^2} + \|h^\delta(t)\|_{L^\infty}\}^2 dt \\ &\leq C e^{2\operatorname{Re}\lambda T^{**}} \{\|\nabla_{x,v}g^\delta(0)\|_{L^2} + \|h^\delta(0)\|_{L^\infty}\}^2 \\ &< C \{e^{\operatorname{Re}\lambda T^\delta} \delta\} C \{e^{\operatorname{Re}\lambda T^{**}} \delta\} \\ &\leq C\theta e^{\operatorname{Re}\lambda T^{**}} \delta. \end{aligned}$$

This is a contradiction to T^{**} in (6.15) when θ is small.

Now that $T^\delta \leq \min\{T^*, T^{**}\}$, we can evaluate $t = T^\delta$ in (6.12) to get

$$\|g^\delta(T^\delta) - \delta e^{-\mathcal{L}T^\delta} \operatorname{Im}R^\delta\|_{L^2} \leq C\theta^2.$$

But

$$\begin{aligned} \|\delta e^{-\mathcal{L}T^\delta} \operatorname{Im}R^\delta\|_{L^2} &\geq \|\delta e^{-\mathcal{L}T^\delta} \operatorname{Im}R\|_{L^2} - \|\delta e^{-\mathcal{L}T^\delta} \operatorname{Im}\{R - R_\delta\}\|_{L^2} \\ &\geq \zeta e^{-\operatorname{Re}\lambda T^\delta} \|\operatorname{Im}R\|_{L^2} - C_A \sqrt{\delta} e^{\lambda T^\delta} \delta r \\ &= \zeta e^{\operatorname{Re}\lambda T^\delta} \delta r - \frac{\zeta}{4} e^{\operatorname{Re}\lambda T^\delta} \delta r = \frac{\zeta}{4} r\theta. \end{aligned}$$

Therefore, for θ sufficiently small,

$$\|g^\delta(T^\delta)\|_{L^2} \geq \frac{\zeta}{4} r\theta - C\theta^2 \geq \frac{\zeta}{8} r\theta.$$

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